

Nuclear astrophysics of the s- and r-process

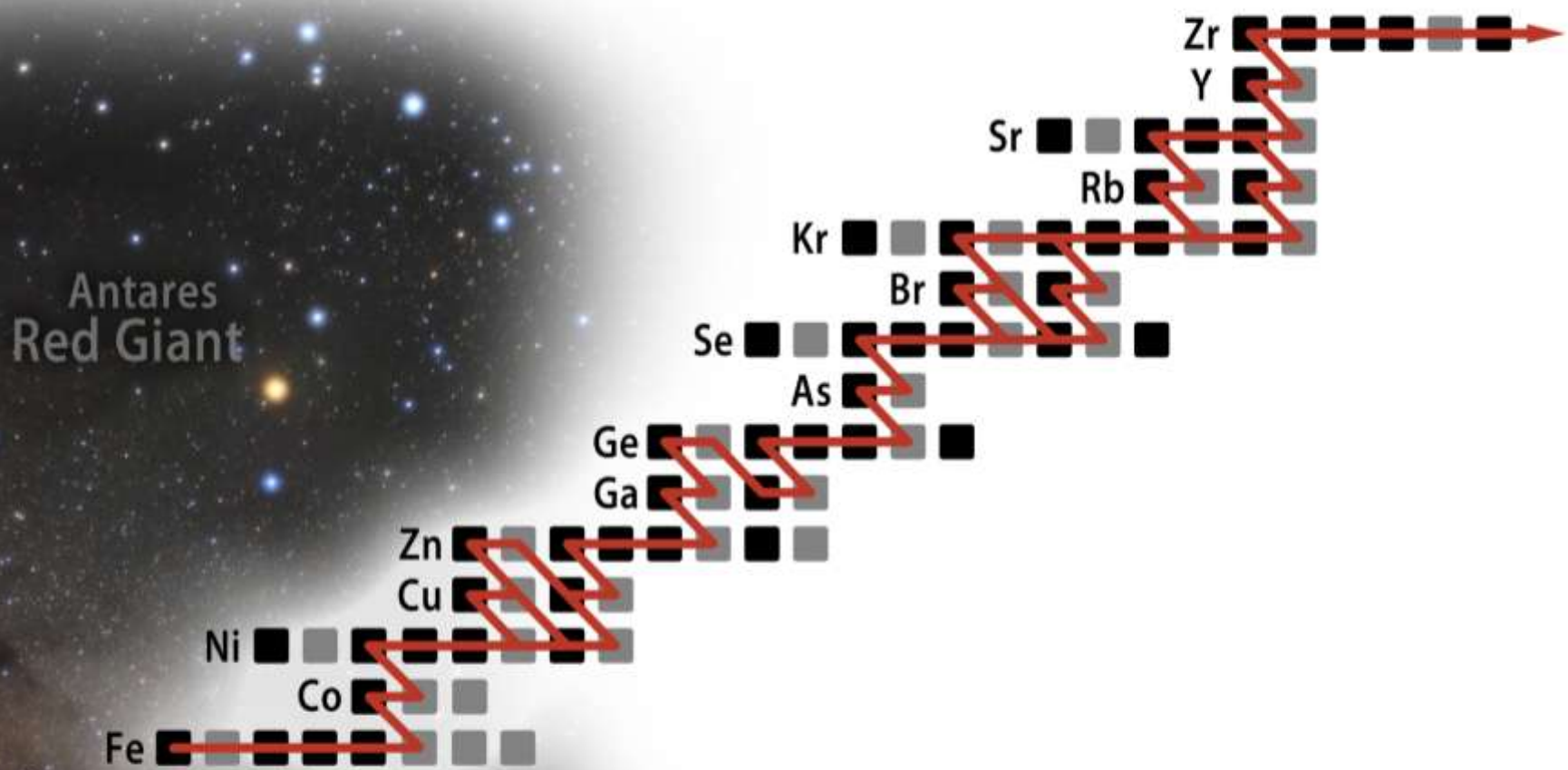
René Reifarth

Goethe University Frankfurt

***Ecole Joliot Curie School
on “Neutrons and Nuclei”***

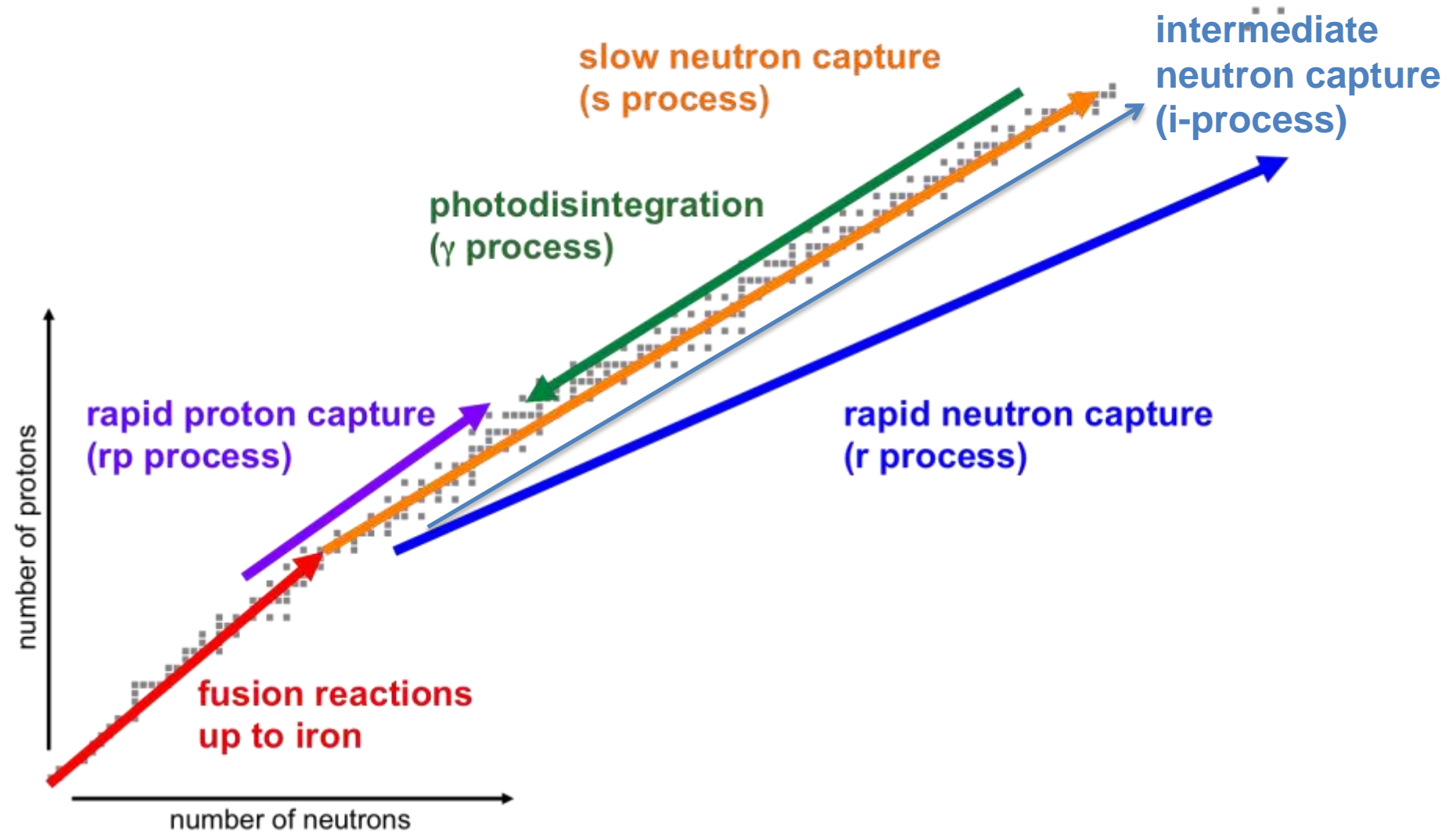
Frejus, France, Sep-28 – Oct-3 2014

Nucleosynthesis – tales from the past

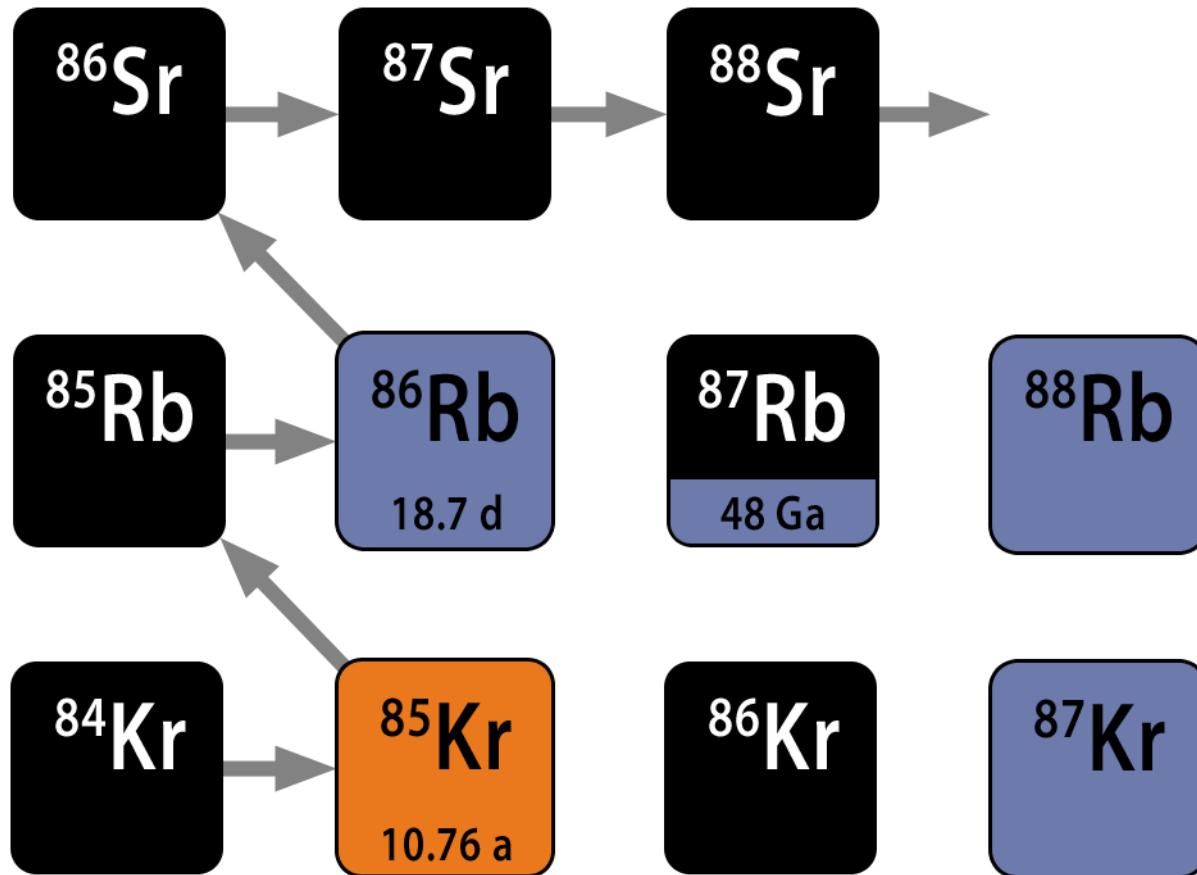




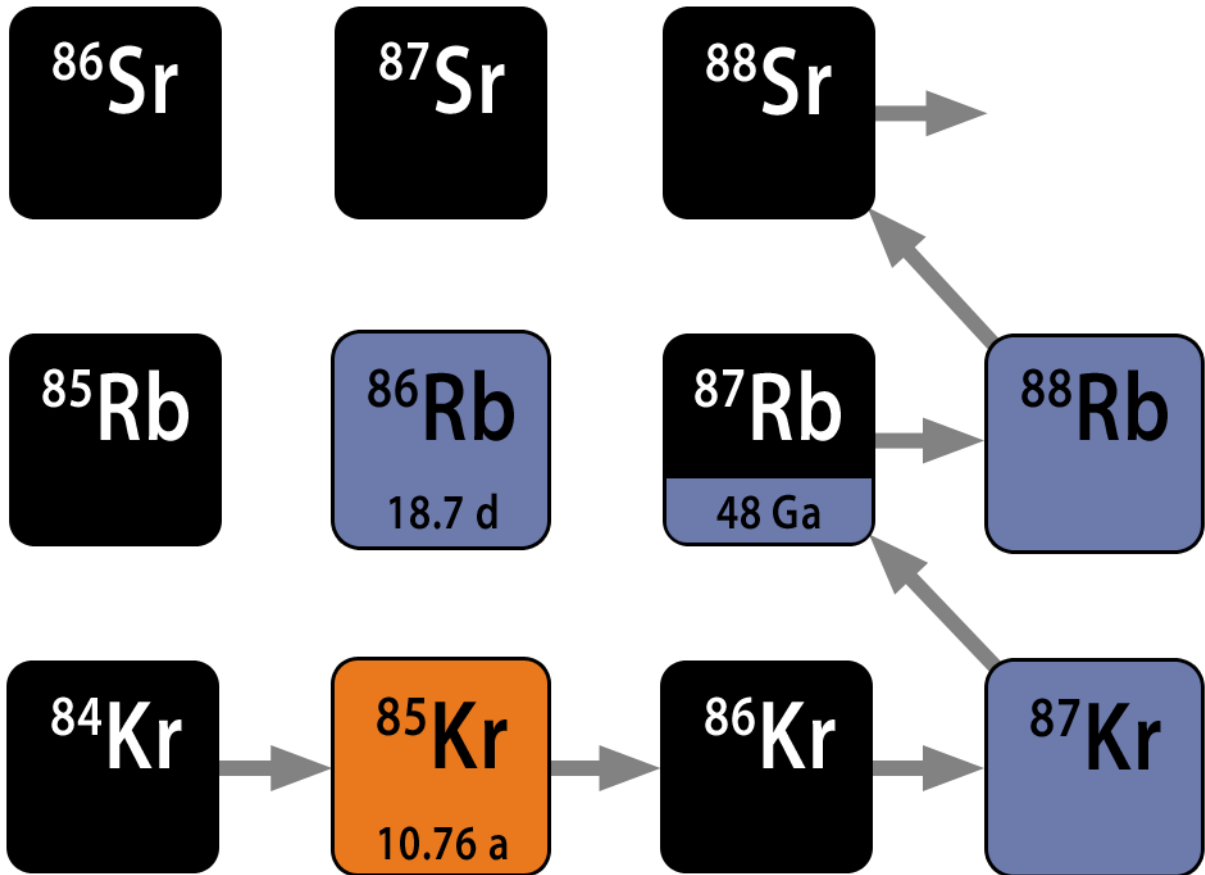
The nucleosynthesis of the elements



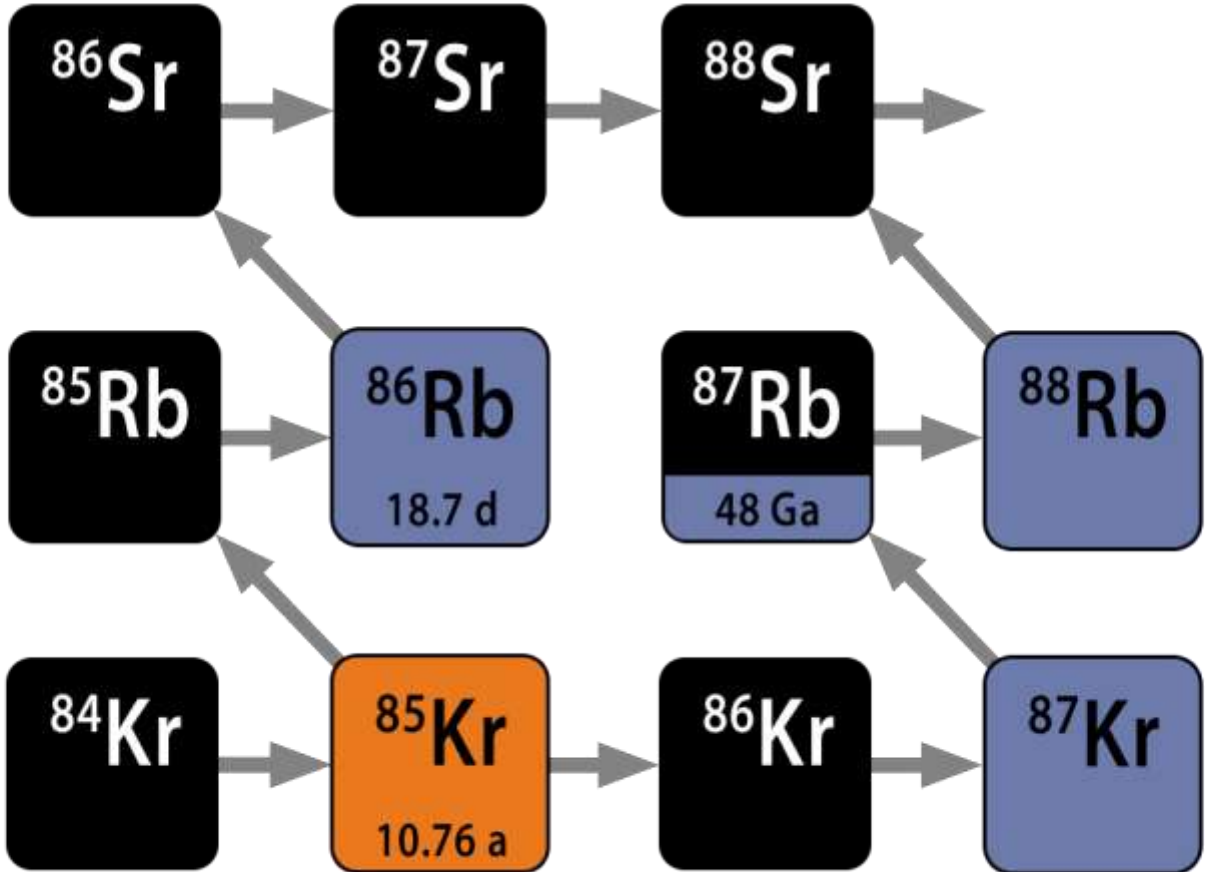
Radioactive isotopes in the s-process



Radioactive isotopes in the s-process

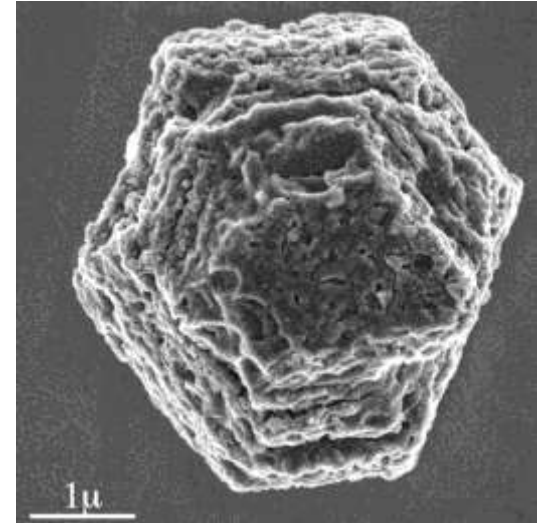
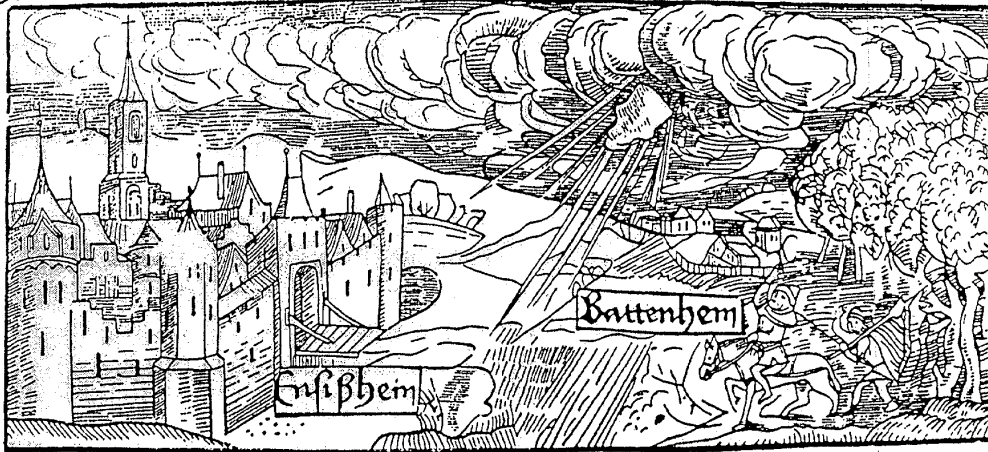


Radioactive isotopes in the s-process

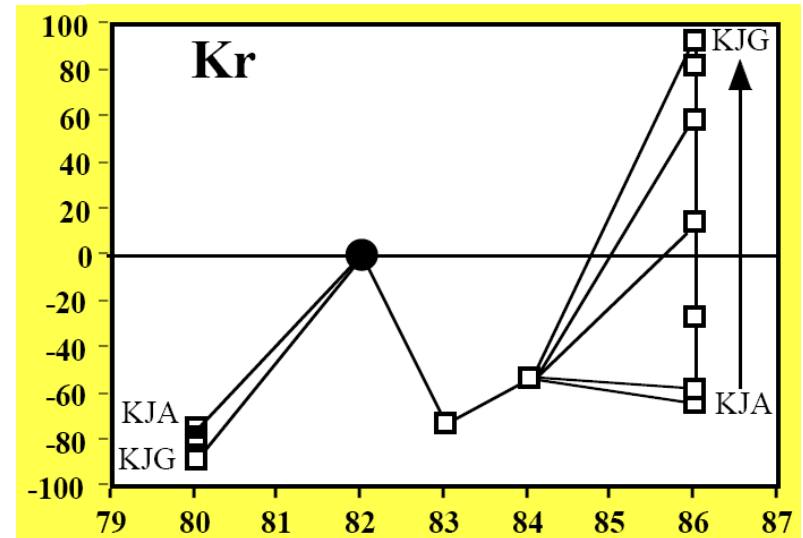


Meteorites – hints from the sky

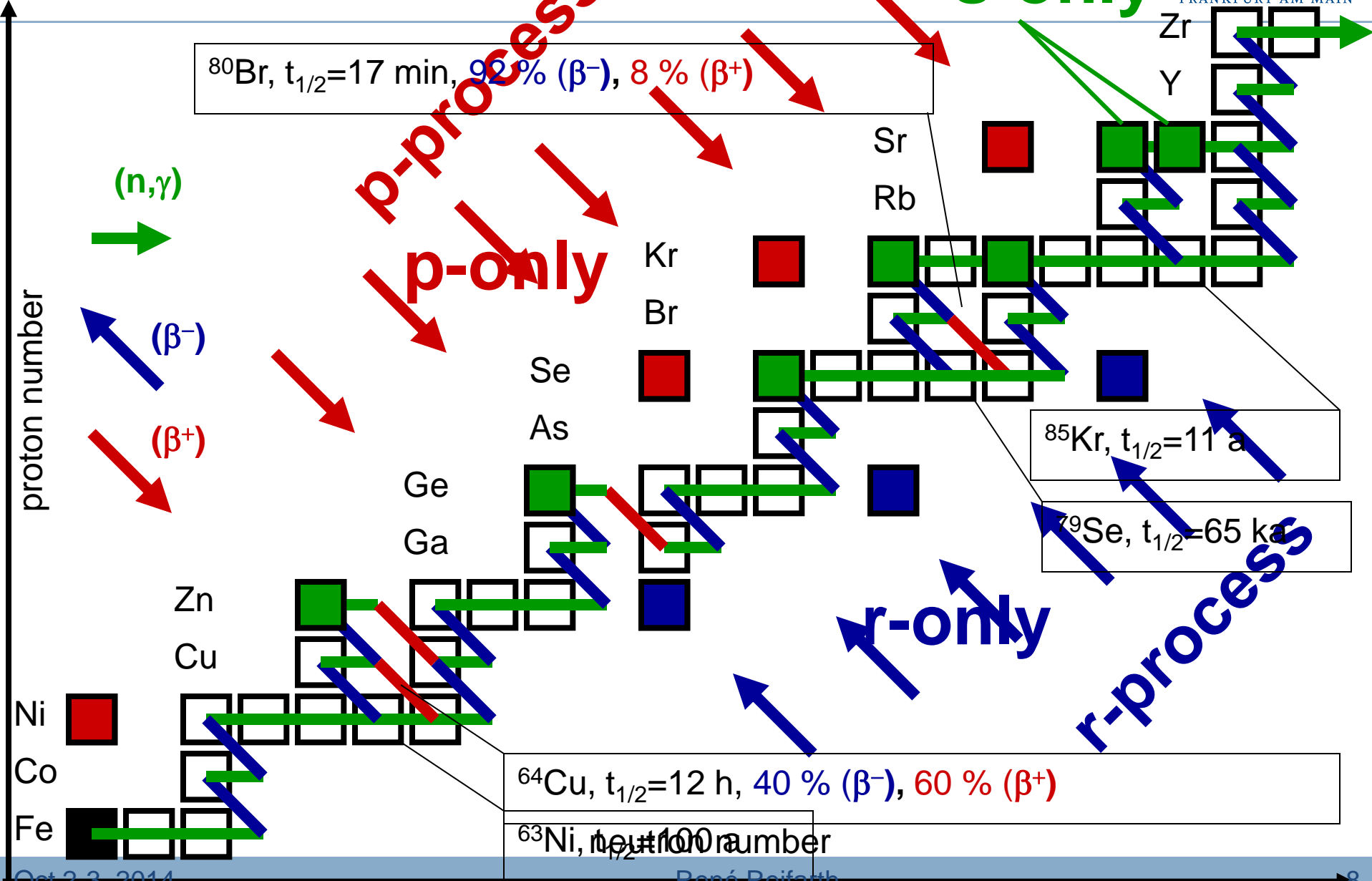
Von dem donnerstein gefallē im xviij. iar: vor Ensisheim.



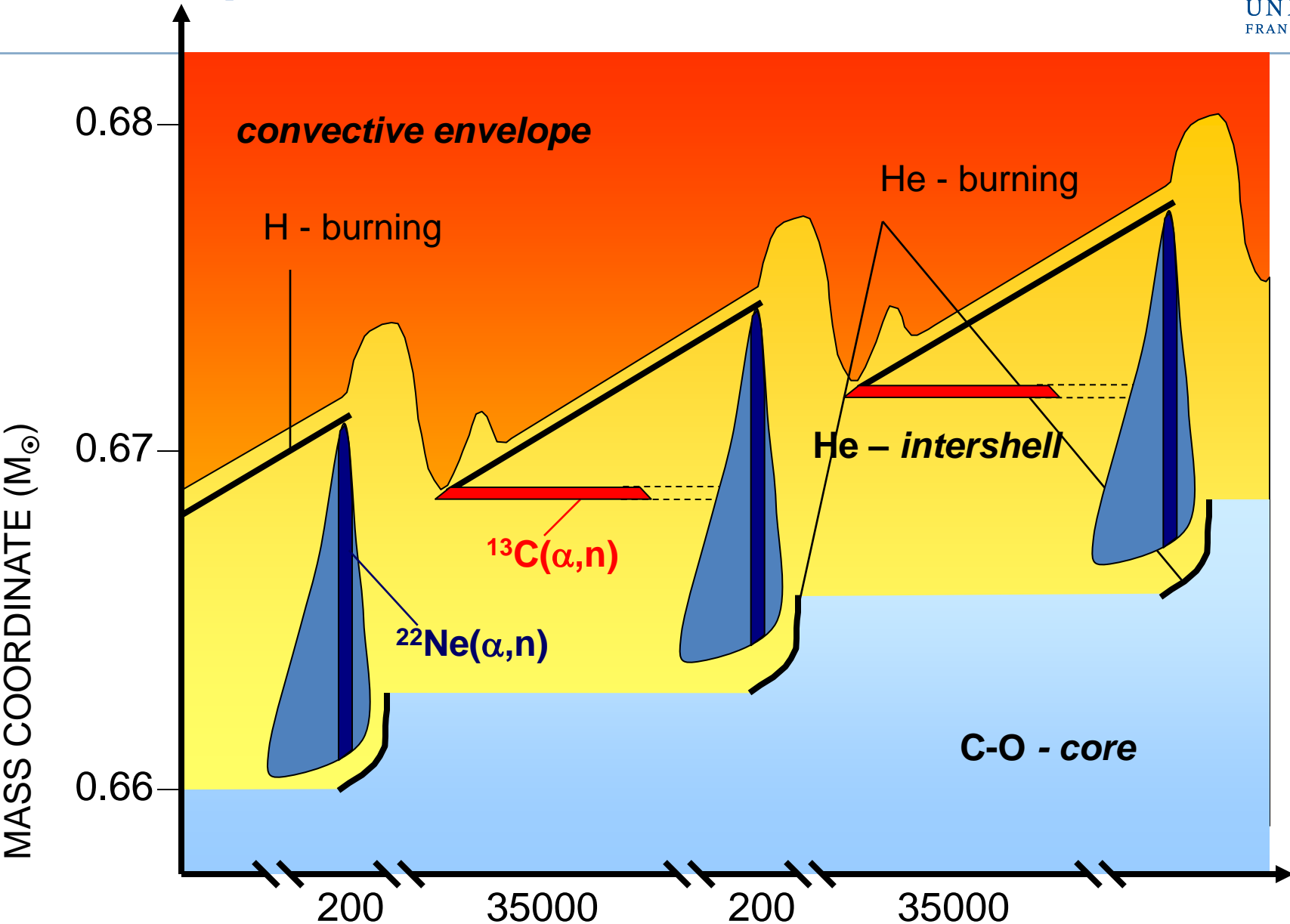
Meteorites contain presolar grains!



the s-process



s-process in AGB stars



s-process nucleosynthesis

Two components were identified and connected to stellar sites:

Main s-process $90 < A < 210$

TP-AGB stars $1-3 M_{\odot}$

shell H-burning
 $0.9 \cdot 10^8$ K

He-flash
 $3-3.5 \cdot 10^8$ K

$kT = 8$ keV

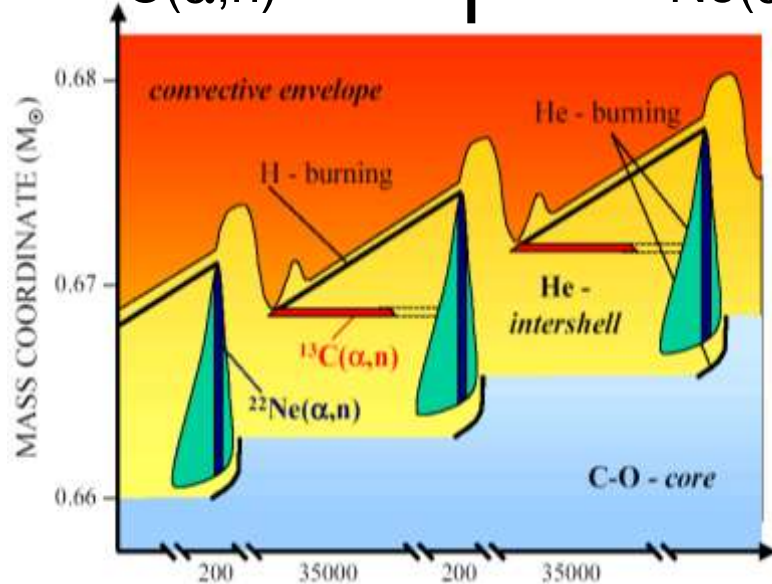
$kT = 25$ keV

10^7-10^8 cm $^{-3}$

$10^{10}-10^{11}$ cm $^{-3}$

$^{13}\text{C}(\alpha, n)$

$^{22}\text{Ne}(\alpha, n)$



Weak s-process $A < 90$

massive stars $> 8 M_{\odot}$

core He-burning
 $3-3.5 \cdot 10^8$ K

shell C-burning
 $\sim 1 \cdot 10^9$ K

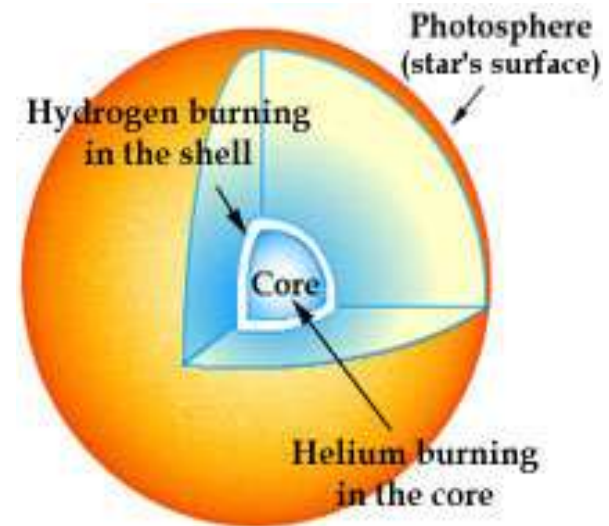
$kT = 25$ keV

$kT = 90$ keV

10^6 cm $^{-3}$

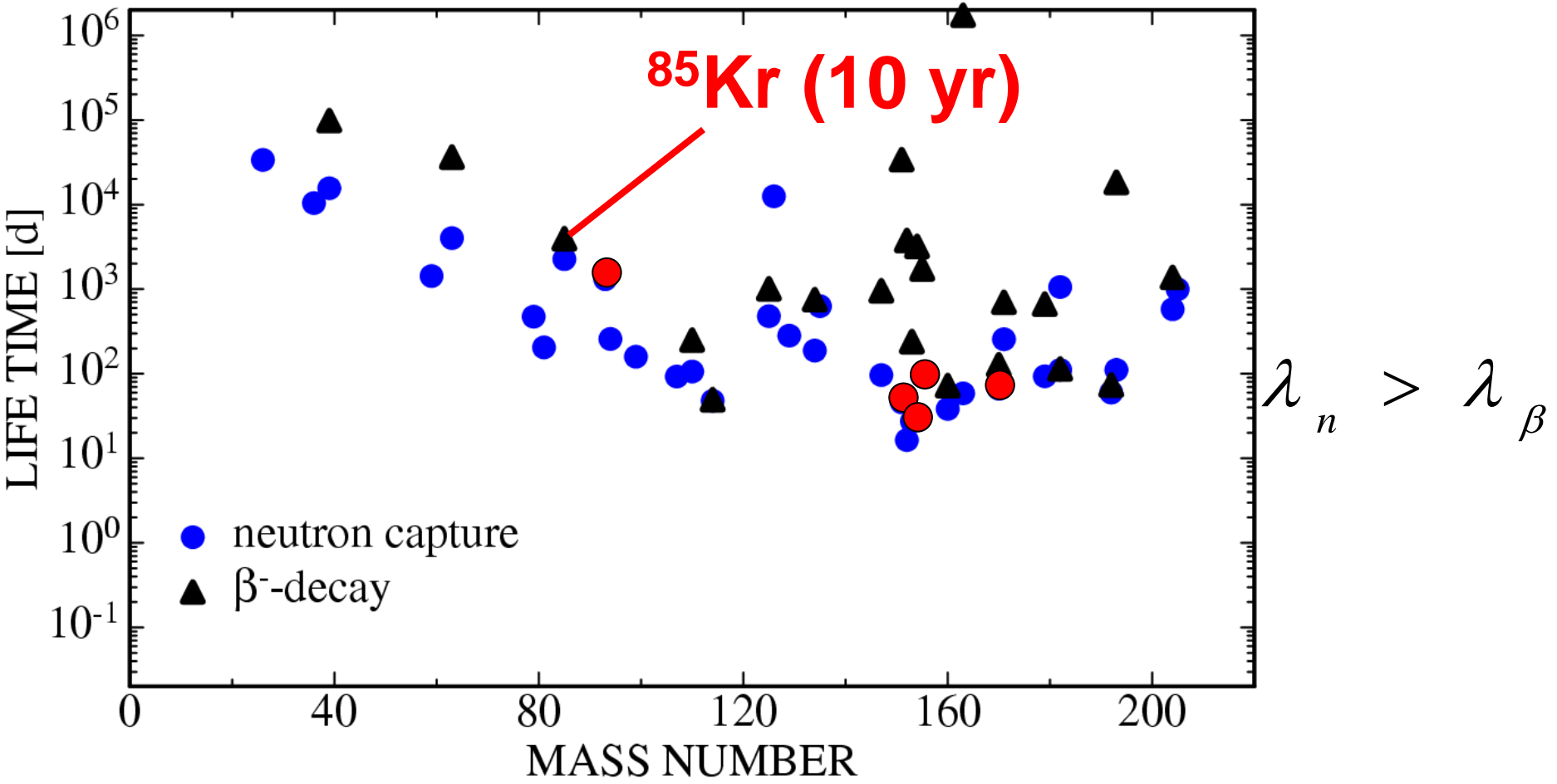
$10^{11}-10^{12}$ cm $^{-3}$

$^{22}\text{Ne}(\alpha, n)$



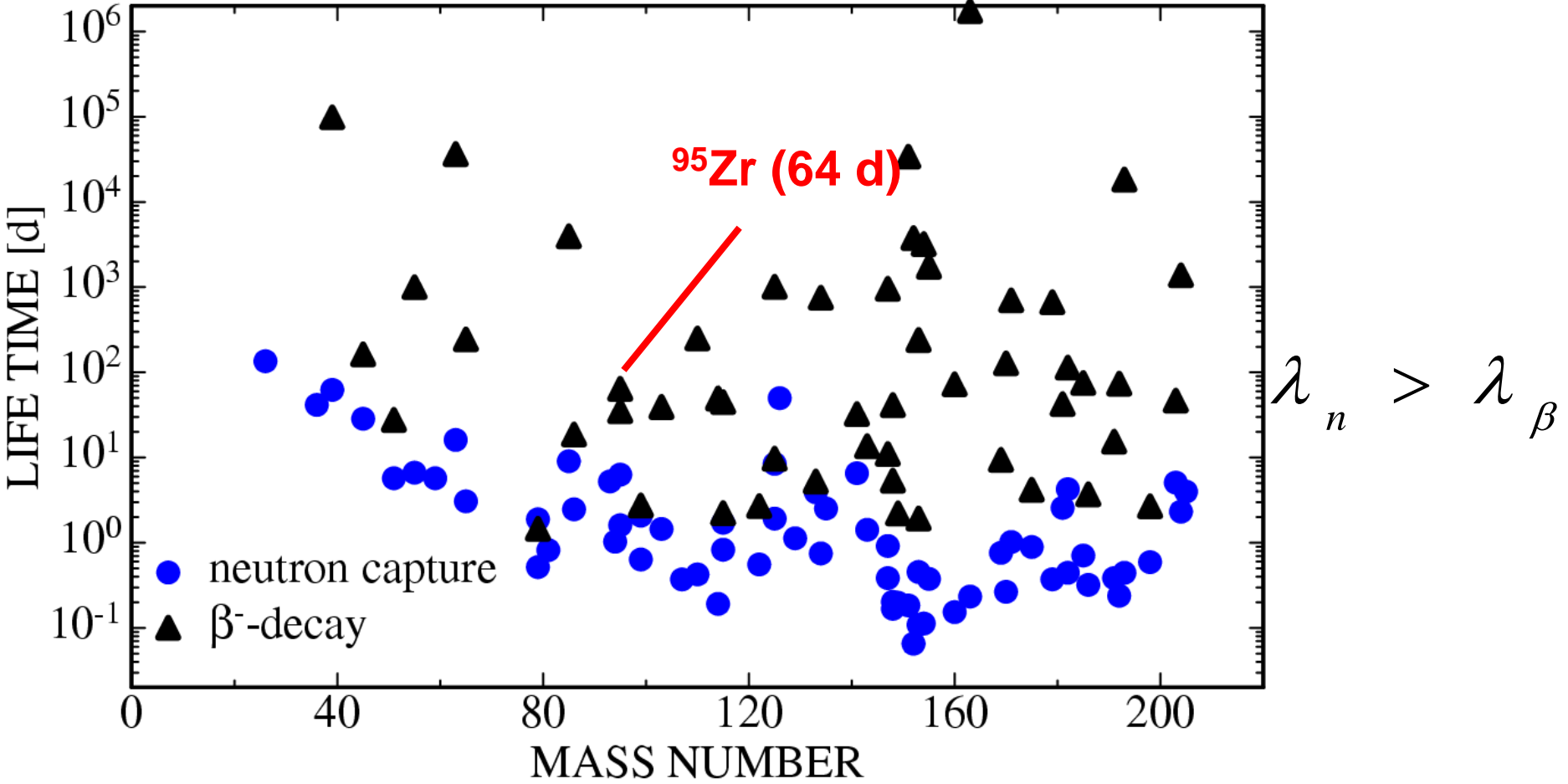
s-process models - classical s-process

Life Times for Unstable Isotopes, $\rho_n = 4 \cdot 10^8 \text{ cm}^{-3}$



s-process models – T-AGB stars, ^{22}Ne phase

Life Times for Unstable Isotopes, $\rho_n = 10^{11} \text{ cm}^{-3}$



Couture & Reifarth, ADNDT, 93 (2007) 807

Nucleosynthesis in the r-process

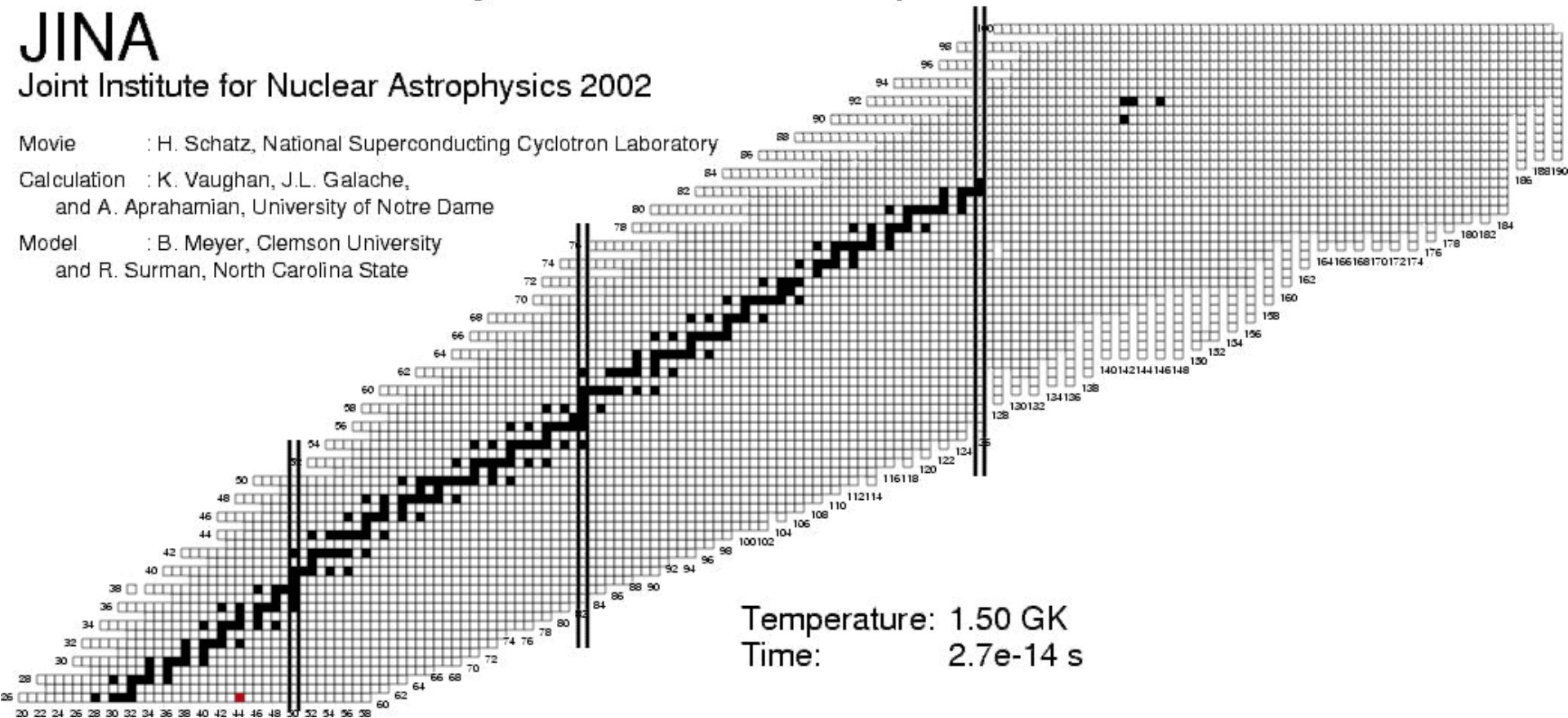
JINA

Joint Institute for Nuclear Astrophysics 2002

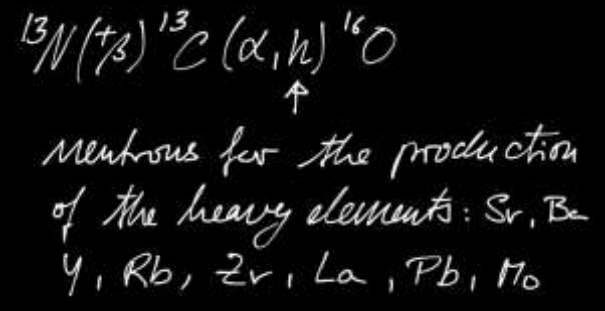
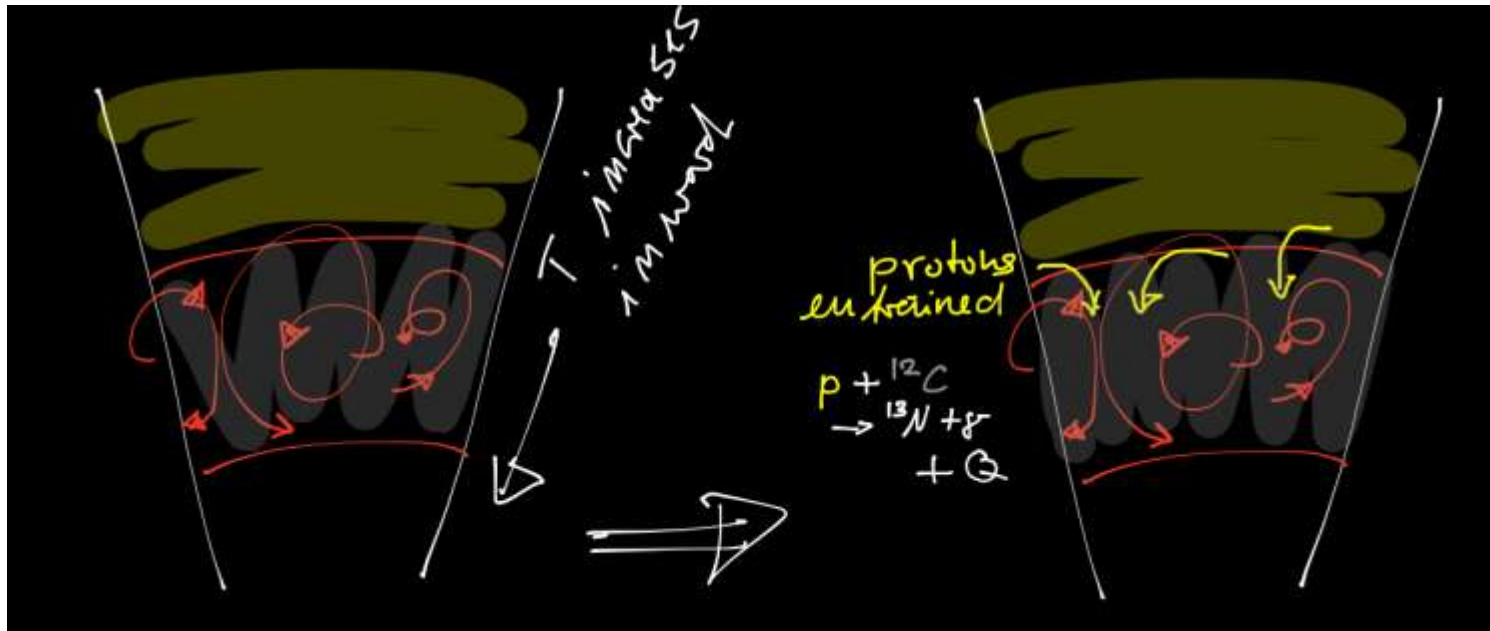
Movie : H. Schatz, National Superconducting Cyclotron Laboratory

Calculation : K. Vaughan, J.L. Galache,
and A. Aprahamian, University of Notre Dame

Model : B. Meyer, Clemson University
and R. Surman, North Carolina State



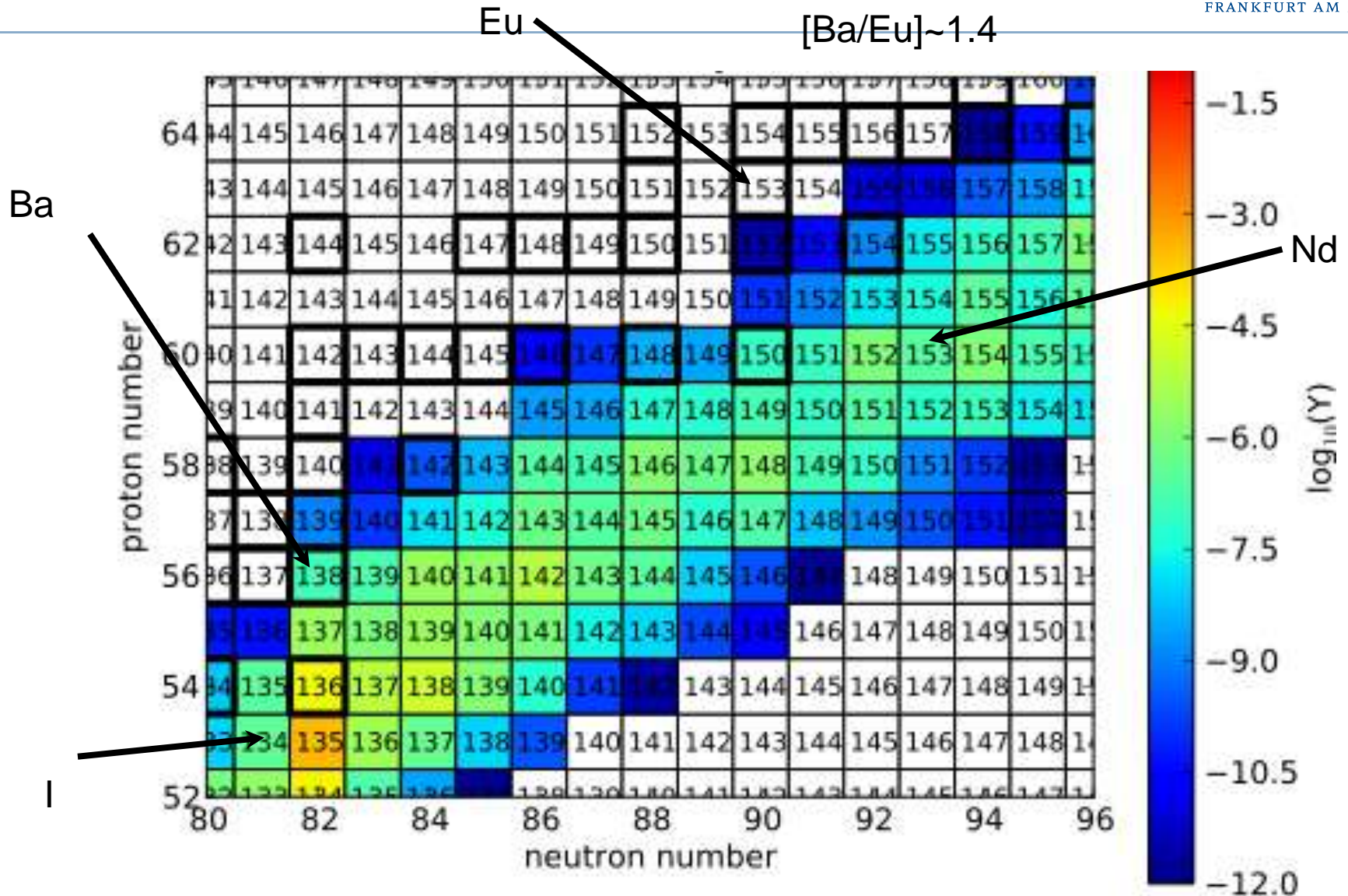
H-entrainment into He-shell flash convection zone



Much higher neutron densities, since ${}^{13}\text{N}$ get mixed deeply into the hot zones. ${}^{13}\text{C}$ gets processed in minutes instead of 1000s of years.

F. Herwig, *The Astrophysical Journal* 727 (2011) 89

The i-process path

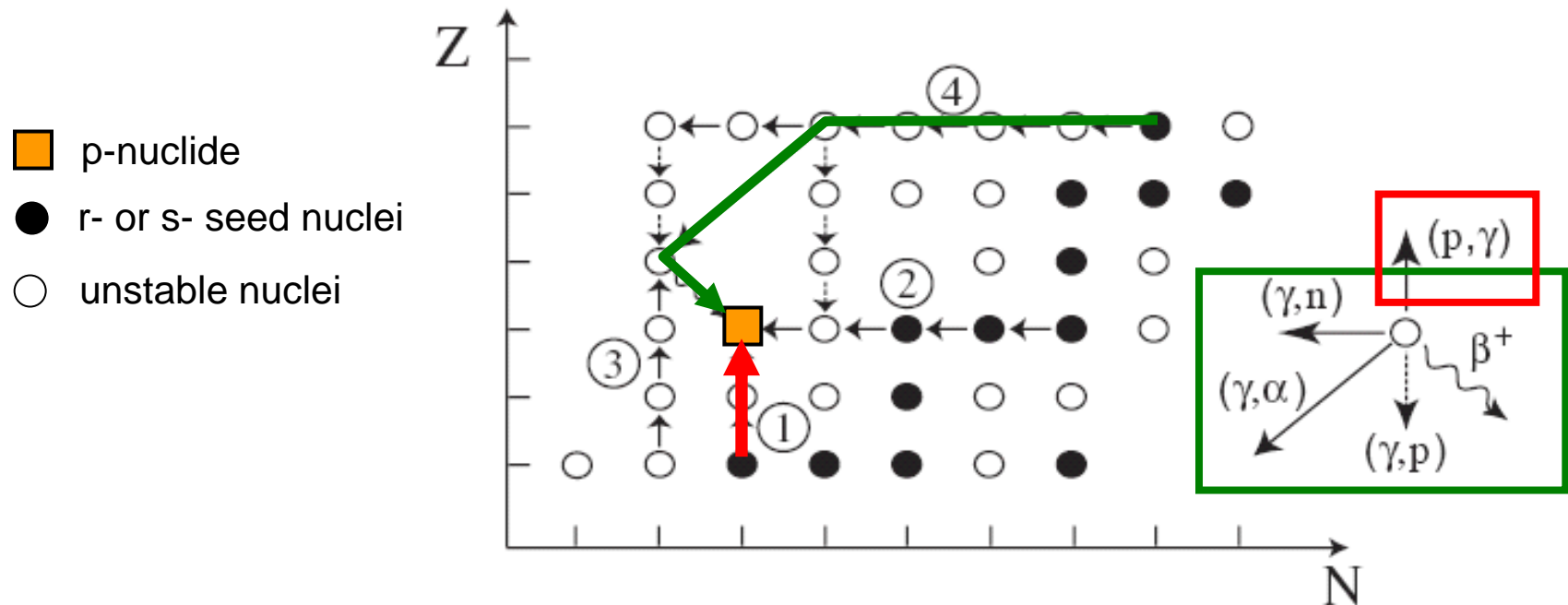


Key reaction in the i-process



the p-process

- 35 stable neutron-deficient isotopes between ^{74}Se and ^{196}Hg
- Dominating reactions: **(p, γ)** for light nuclei;
(γ ,n), (γ ,p), (γ , α) and **β^+** decays for heavier nuclei
- Temperatures of $2\text{-}3 \times 10^9$ K during time scales of a few seconds are required (type II supernovae explosions)



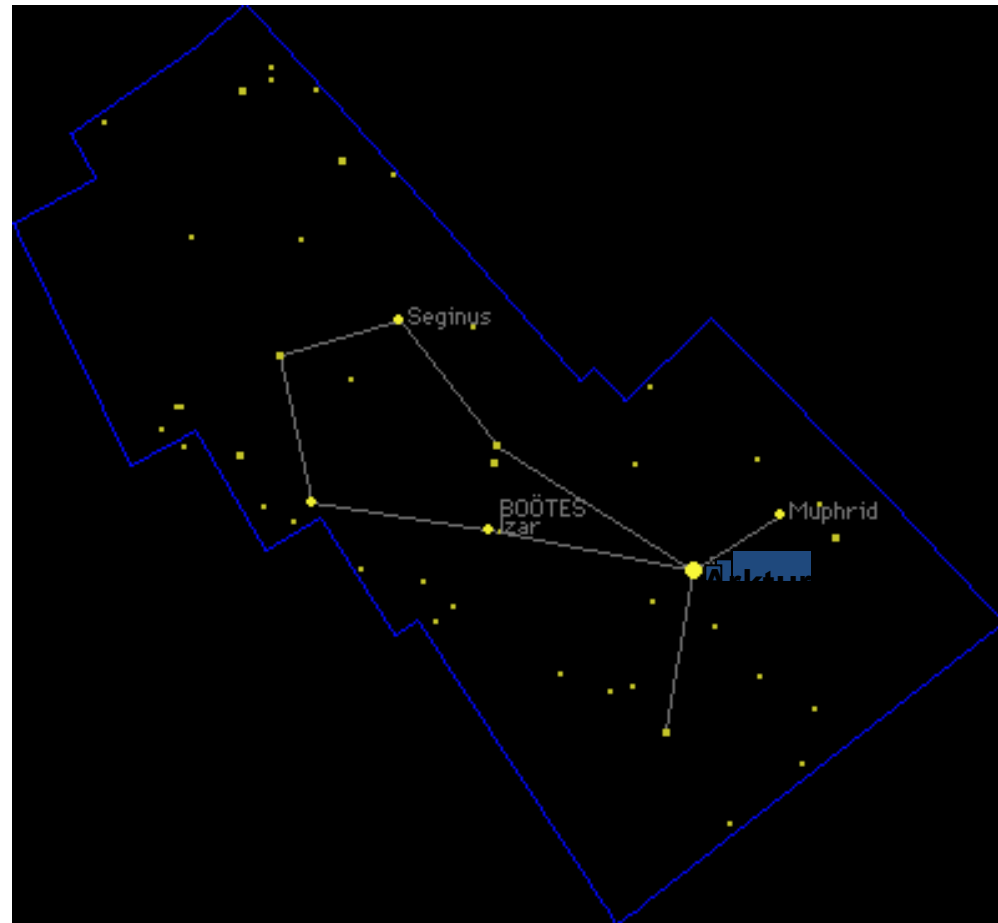
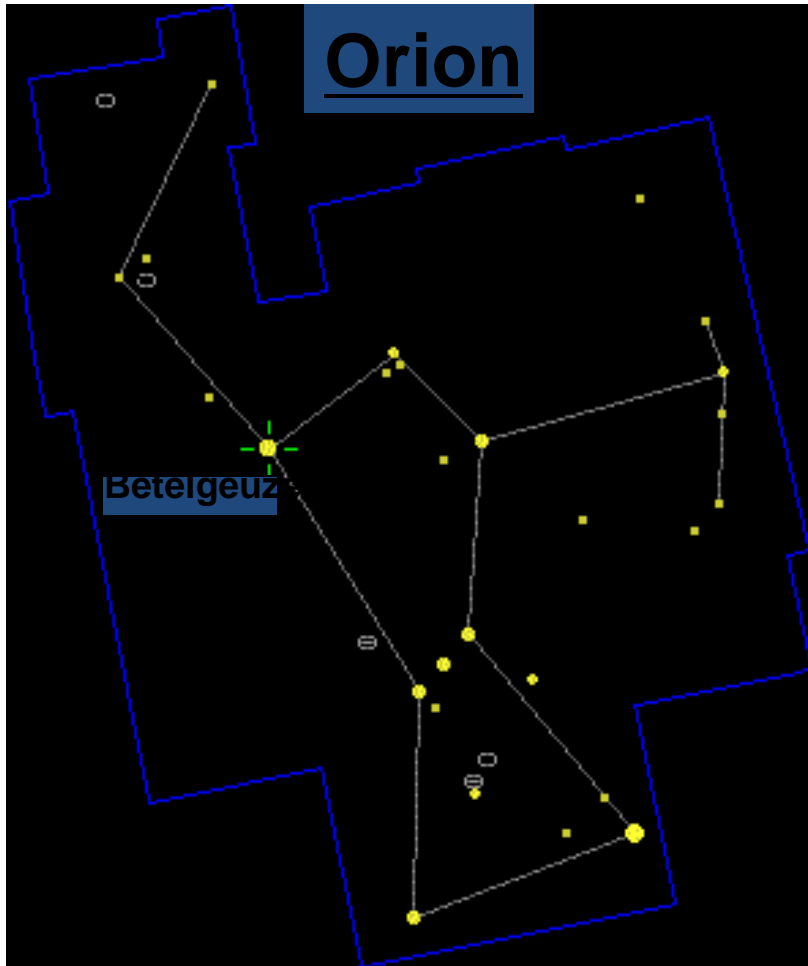
(n,γ) reactions in the p-process

- very high temperatures
- γ-induced reactions
- result: free neutrons and neutron-deficient material

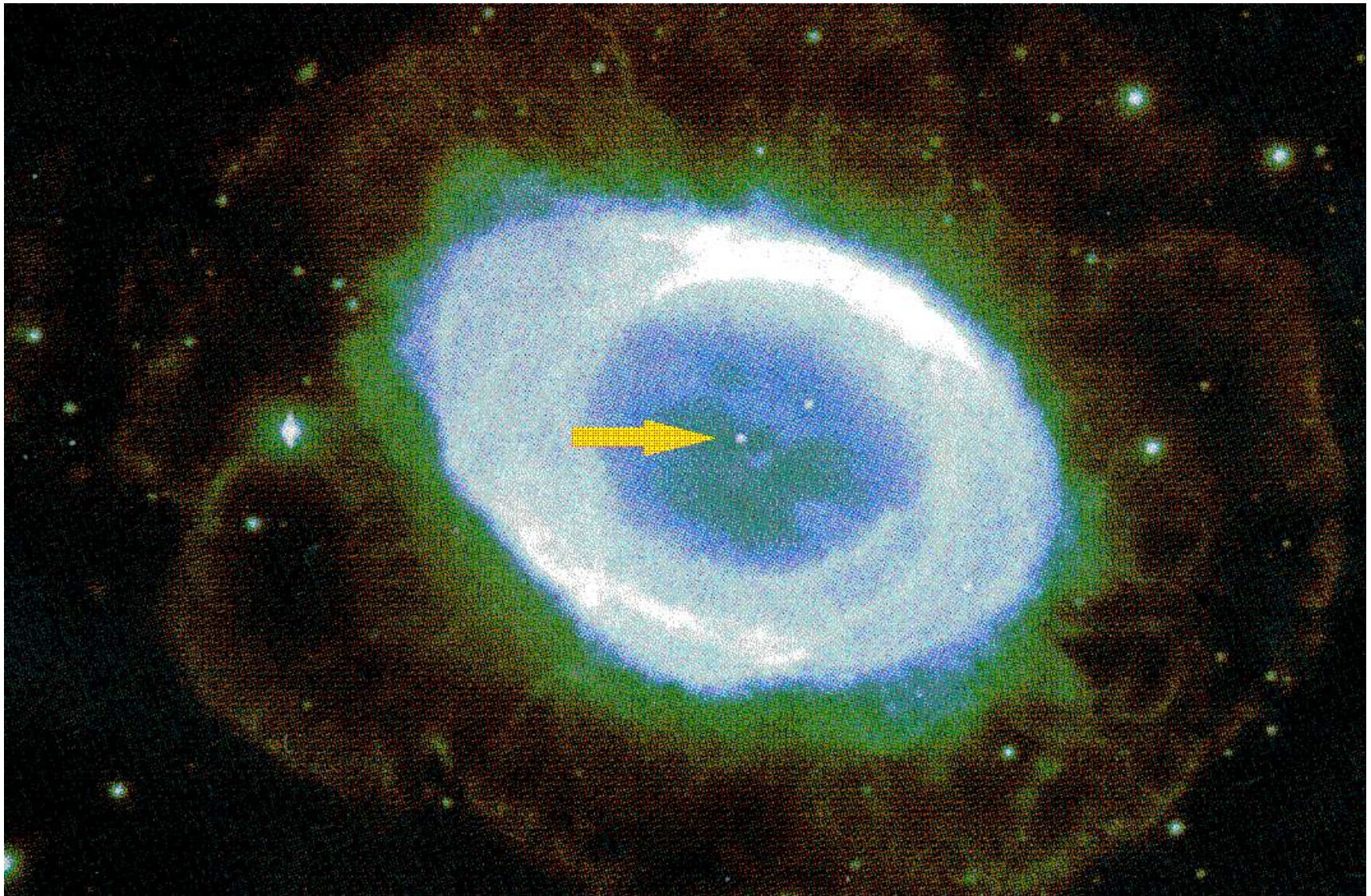


Arnould & Goriely, Physics Reports 384 (2003) 1–84

Red Giants – easy to spot



Red Giants become White Dwarfs

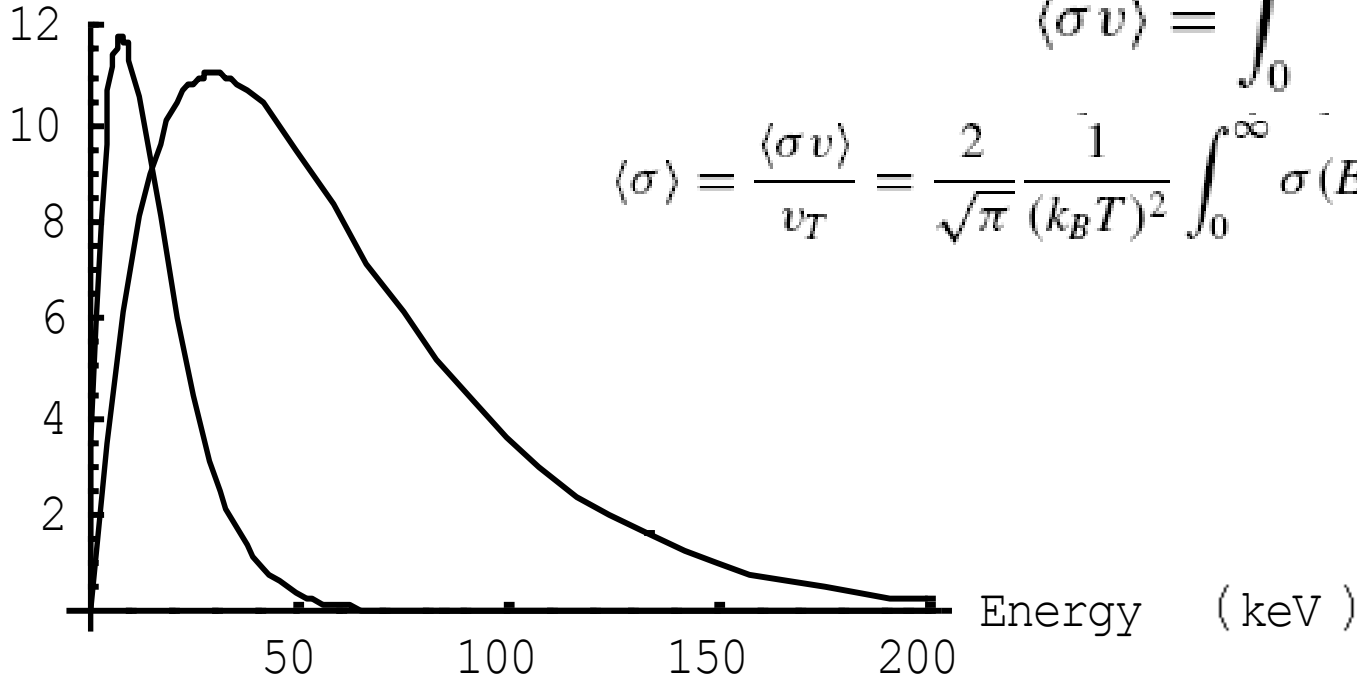


Ring nebula illuminated by the White Dwarf in the center.

What's needed?

Neutron induced Reaction rates (1-200 keV)

Neutrons



$$\langle \sigma v \rangle = \int_0^{\infty} \phi(v) \sigma(v) v dv$$

$$\langle \sigma \rangle = \frac{\langle \sigma v \rangle}{v_T} = \frac{2}{\sqrt{\pi}} \frac{1}{(k_B T)^2} \int_0^{\infty} \sigma(E) E \exp\left(-\frac{E}{k_B T}\right) dE.$$

I 121 2,12 h	I 122 3,6 m	I 123 13,2 h	I 124 4,15 d	I 125 59,41 d	I 126 13,11 d	I 127 100	I 128 25,0 m	I 129 1,57 · 10 ⁷ a	I 130 9,0 m	I 131 12,36 h	I 131 8,02 d	I 132 83,6 m	I 132 2,30 h	I 133 9 s	I 133 20,0 h	I 134 3,5 m	I 134 52,0 m	I 135 6,61 h
ε; β ⁺ 1,1... γ 212... g	β ⁺ 3,1... γ 564...	ε; no β ⁺ γ 159... g	ε; β ⁺ 2,1... γ 603; 1691; 723...	ε; γ 35; e ⁻ g σ 900	ε; β ⁻ 0,9; 1,3... β ⁺ 1,1... γ 389; 666... σ ~ 10000	ε 6,15	ε; β ⁻ 2,1... β ⁺ ... γ 443; 527... σ 22	β ⁻ 0,2 γ 40 e ⁻ ; g σ 20,7 + 10,3	β ⁻ 1,0; 1,8... e ⁻ β ⁻ 2,5... γ 536... σ 18	β ⁻ 1,0; 1,8... e ⁻ β ⁻ 2,5... γ 536... σ 18	β ⁻ 0,6; 0,8... γ 364; 637; 284...; g σ ~ 0,7	β ⁻ 2,1... γ 98 β ⁻ 1,5... γ 568; 773; 773; 600; 175...	β ⁻ 2,1... γ 568; 773; 955; 523...	β ⁻ 1,2; 1,5... γ 913; 647; 73	β ⁻ 1,2; 44 β ⁻ 2,5 γ 847; 884; 234	β ⁻ 1,3; 2,4... γ 847; 884...	β ⁻ 1,5; 2,2... γ 1260; 1132; 1678; 1458... g; m	

Activation Method

$^{14}\text{C}(n,\gamma)^{15}\text{C}$ reaction

detected via

$^{15}\text{C}(\beta^-)^{15}\text{N}$ decay

($t_{1/2}=2.5$ s)

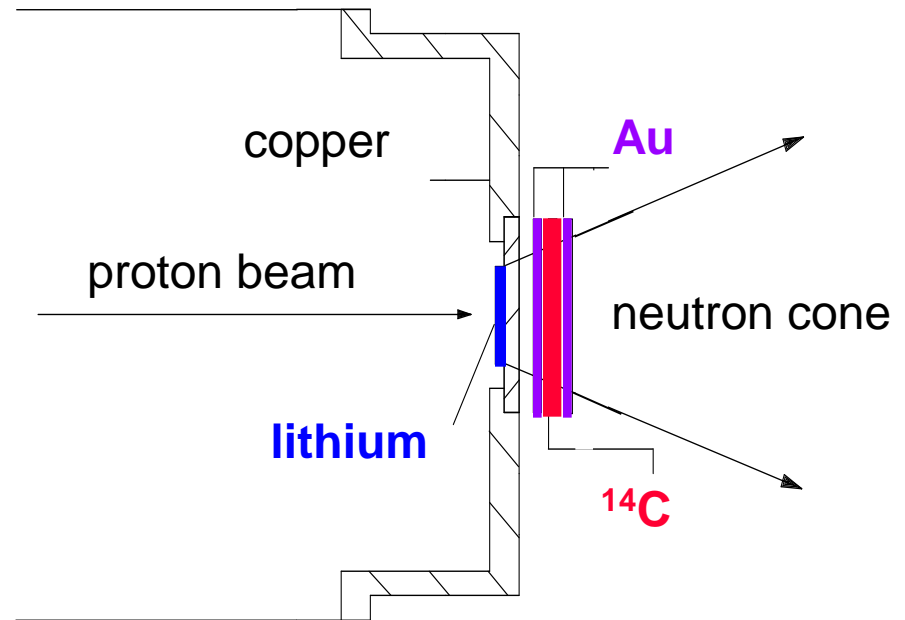
^{14}C sample irradiated for 10 s, then activity counted for 10 s („cyclic activation“)

Determination of
neutron flux via

$^{197}\text{Au}(n,\gamma)^{198}\text{Au}$

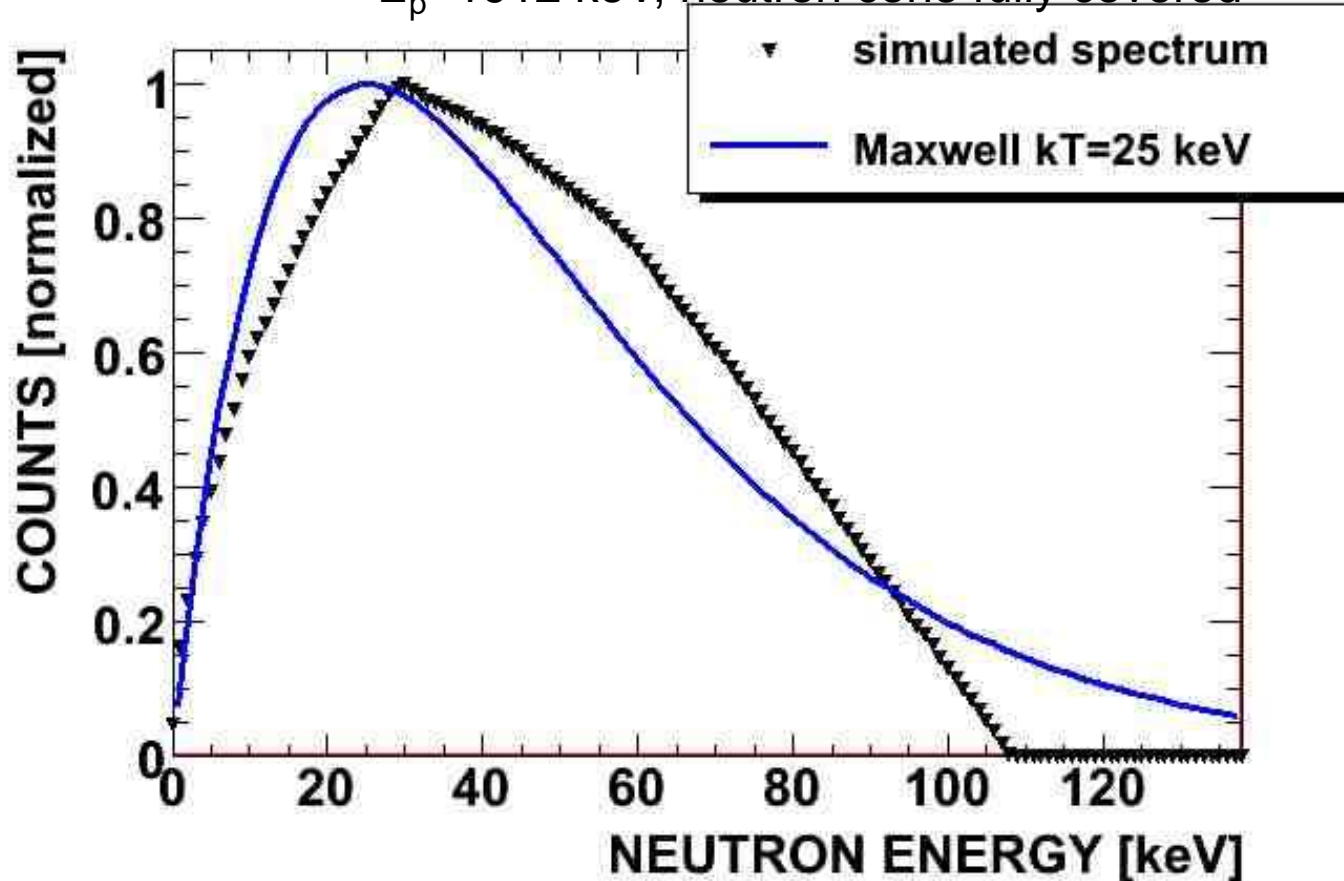
Neutron source:

$^7\text{Li}(p,n)^7\text{Be}$



A standard neutron spectrum – working horse!

$E_p = 1912$ keV, neutron cone fully covered



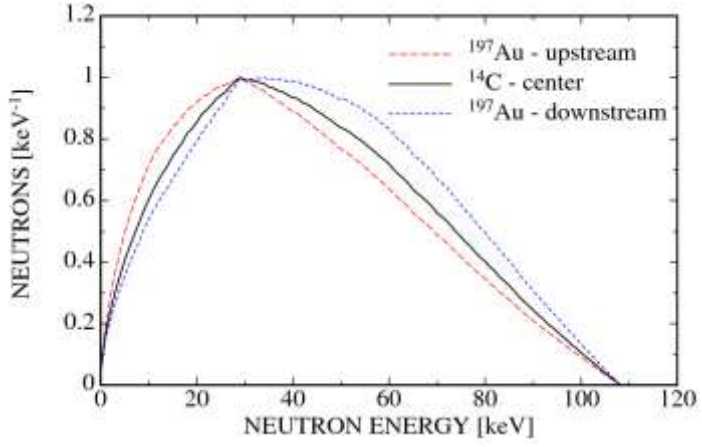
Quasi-Maxwellian averaged
distribution:

$$kT = 25 \text{ keV}$$

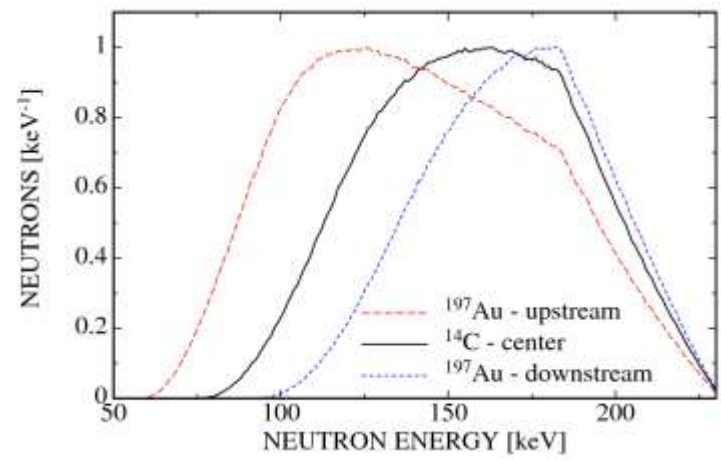
$$E_{max} = 110 \text{ keV}$$

Other neutron spectra

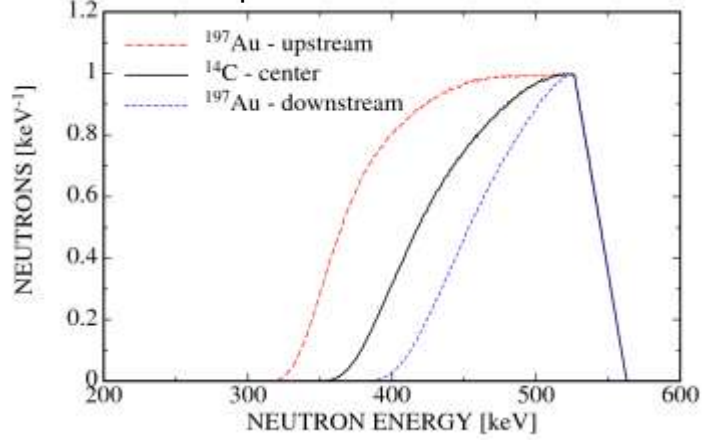
$E_p=1912$ keV



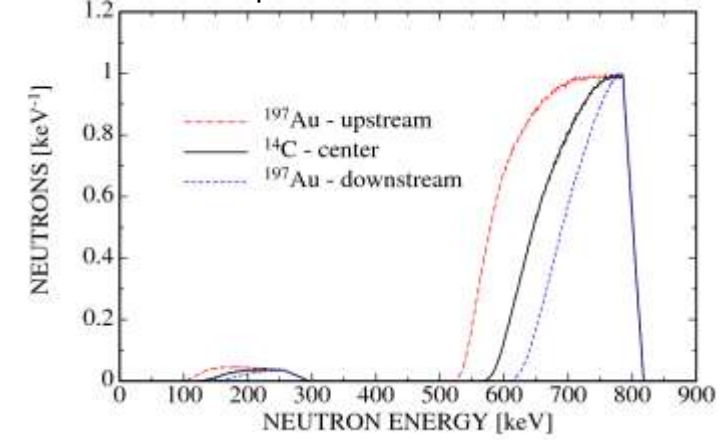
$E_p=2000$ keV



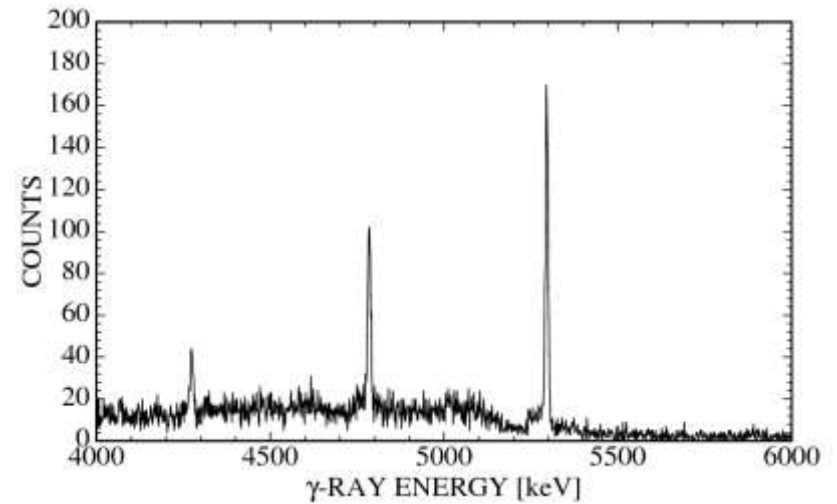
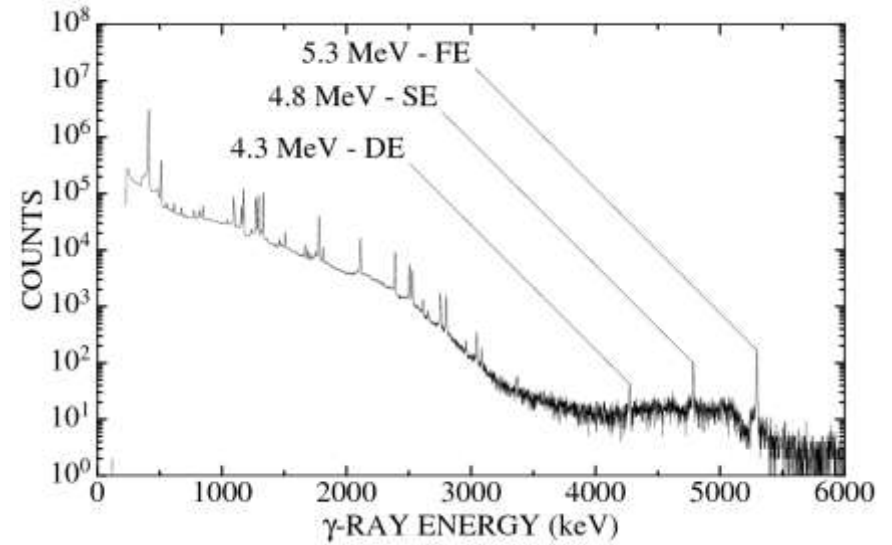
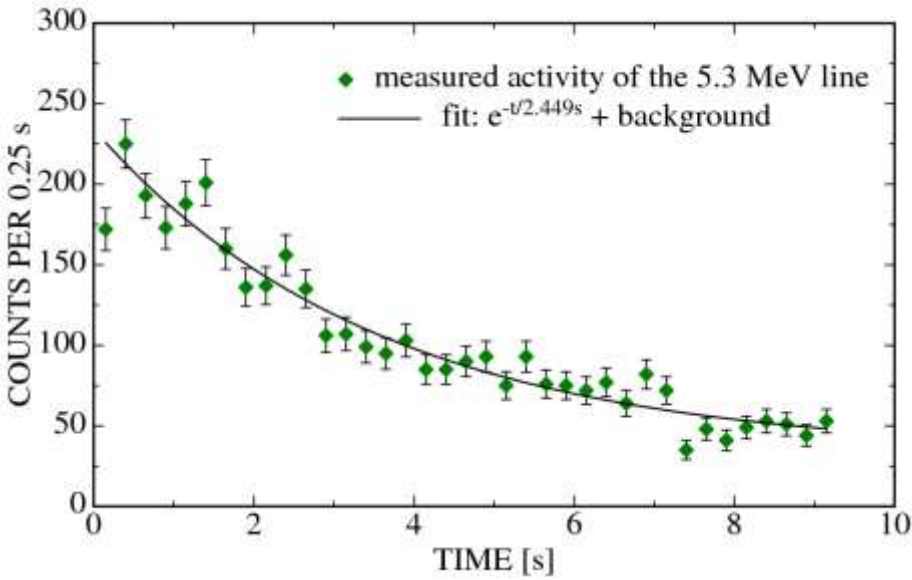
$E_p=2290$ keV



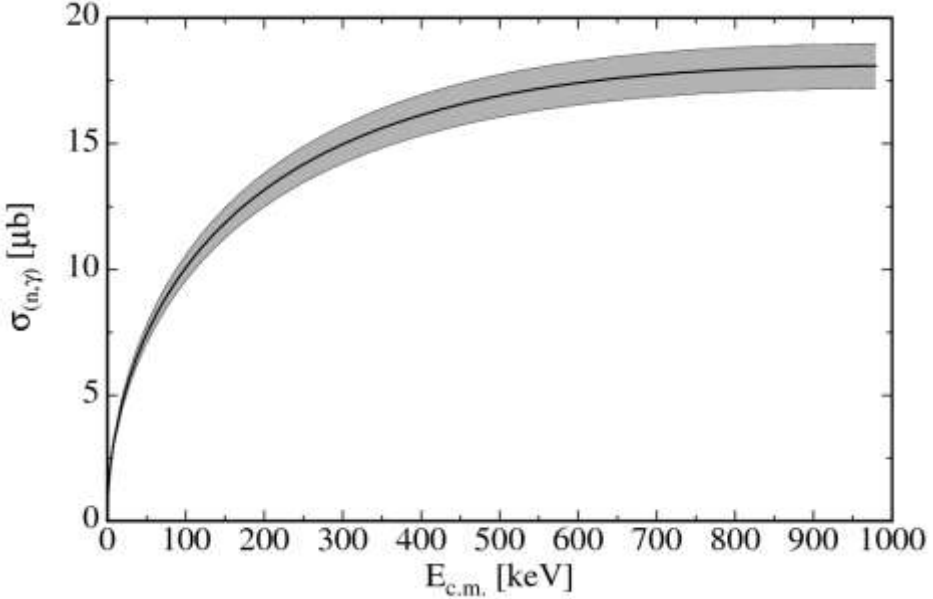
$E_p=2530$ keV



^{15}C – γ -spectra

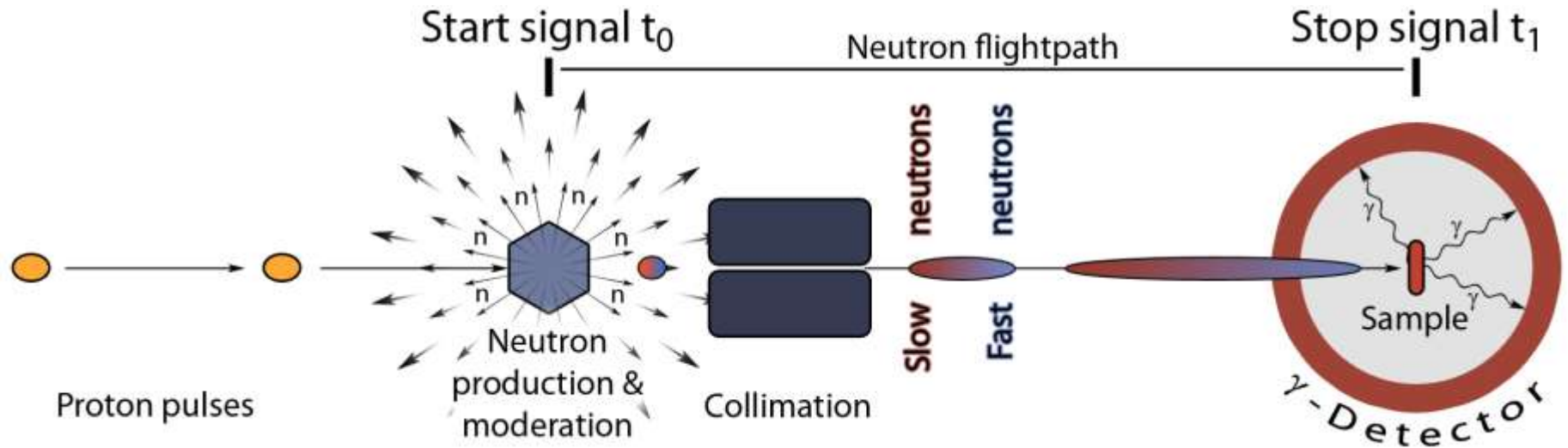


Description and Deconvolution



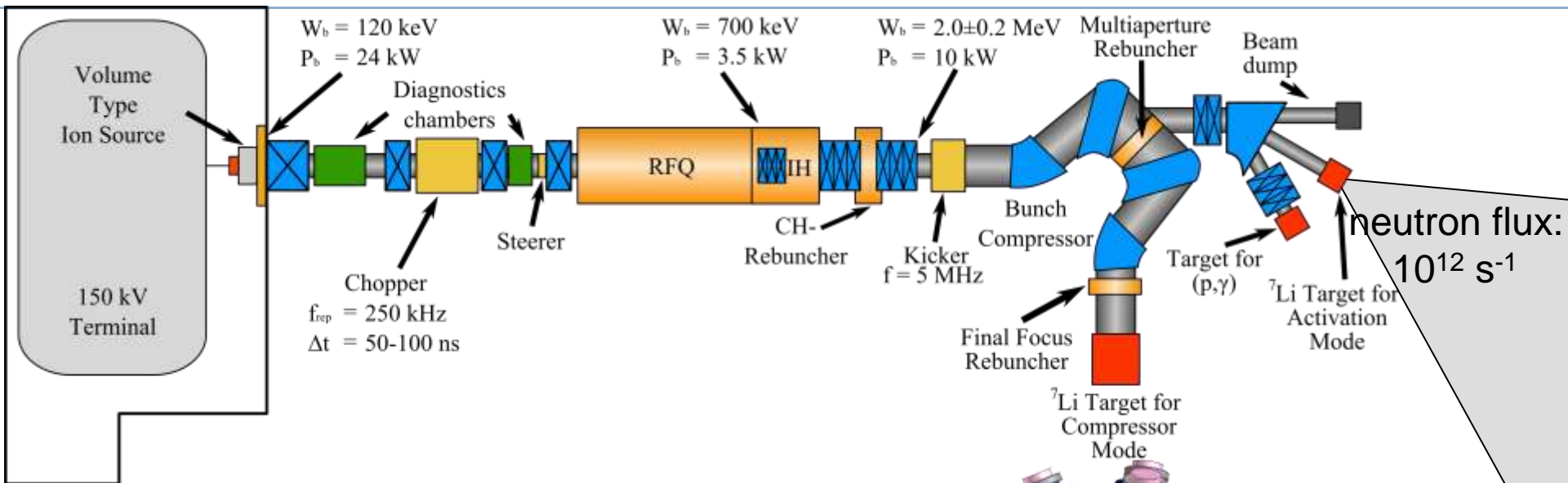
- p-wave capture
- good agreement with exp. data

keV	Exp. [μb]		Theo. [μb]		Theo/Exp	
23	7.1	5	6.5	0.4	0.92	0.08
150	10.7	1.2	11.7	0.6	1.09	0.12
500	17.0	1.5	16.5	0.8	0.97	0.10
800	15.8	1.6	17.5	0.9	1.11	0.11

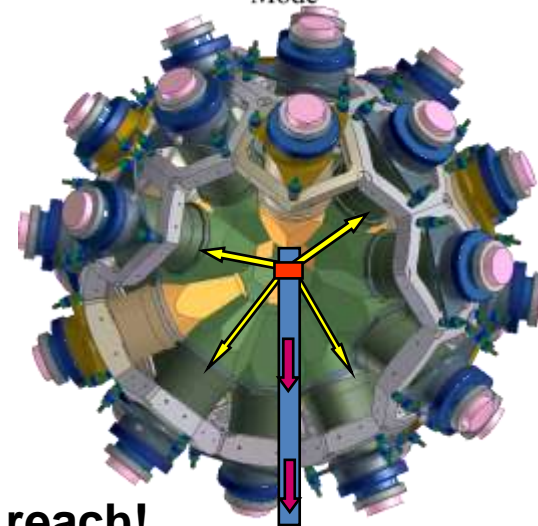


- the TOF-technique is the only generally applicable method to determine energy-dependent neutron capture cross sections
- beam pulsing & distance to the neutron production site significantly reduce the number of neutrons available on the sample

The Frankfurt neutron source at the Stern-Gerlach-Zentrum (FRANZ)



2 mA proton beam (8 A peak current)
 250 kHz
 < 1 ns pulse width
 neutron flux at 1 m: $10^7 \text{ s}^{-1} \text{ cm}^{-2}$
 neutron flux at 0.1 m: $10^9 \text{ s}^{-1} \text{ cm}^{-2}$

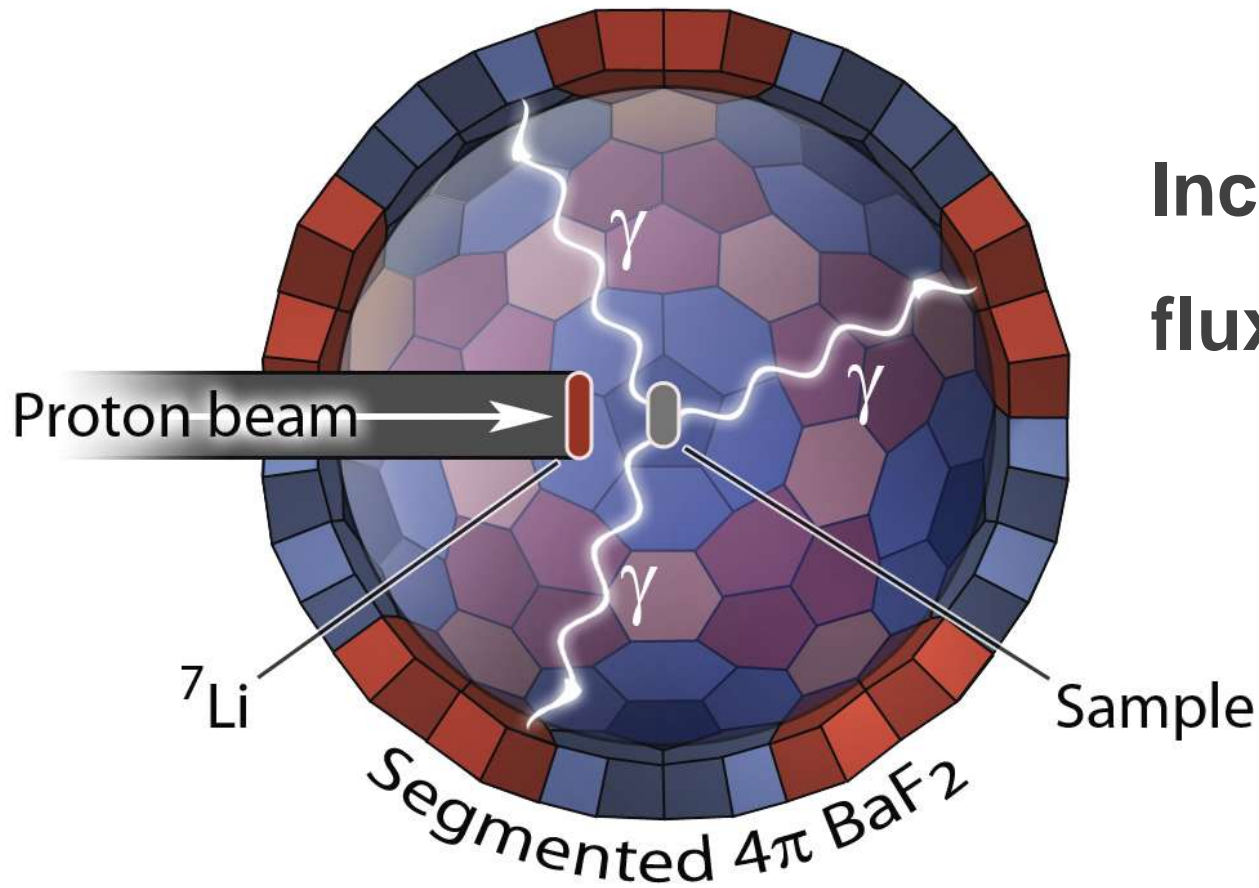


neutron flux:
 10^{12} s^{-1}

Isotopes with half-lives down to months are in reach!

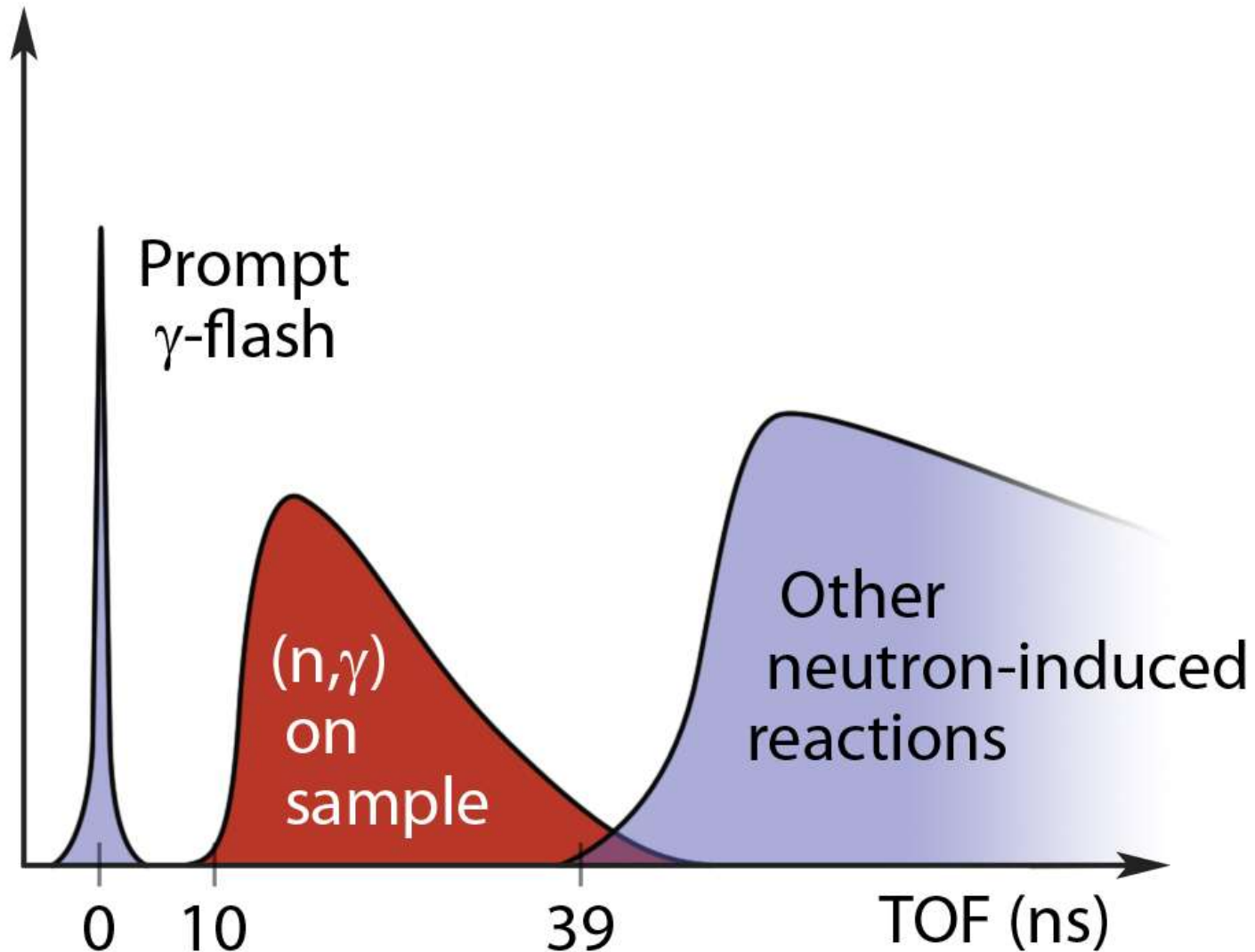
Reifarth et al. PASA 26 (2009) 26, 255–258

NAUTILUS – Neutron capture with short flightpath



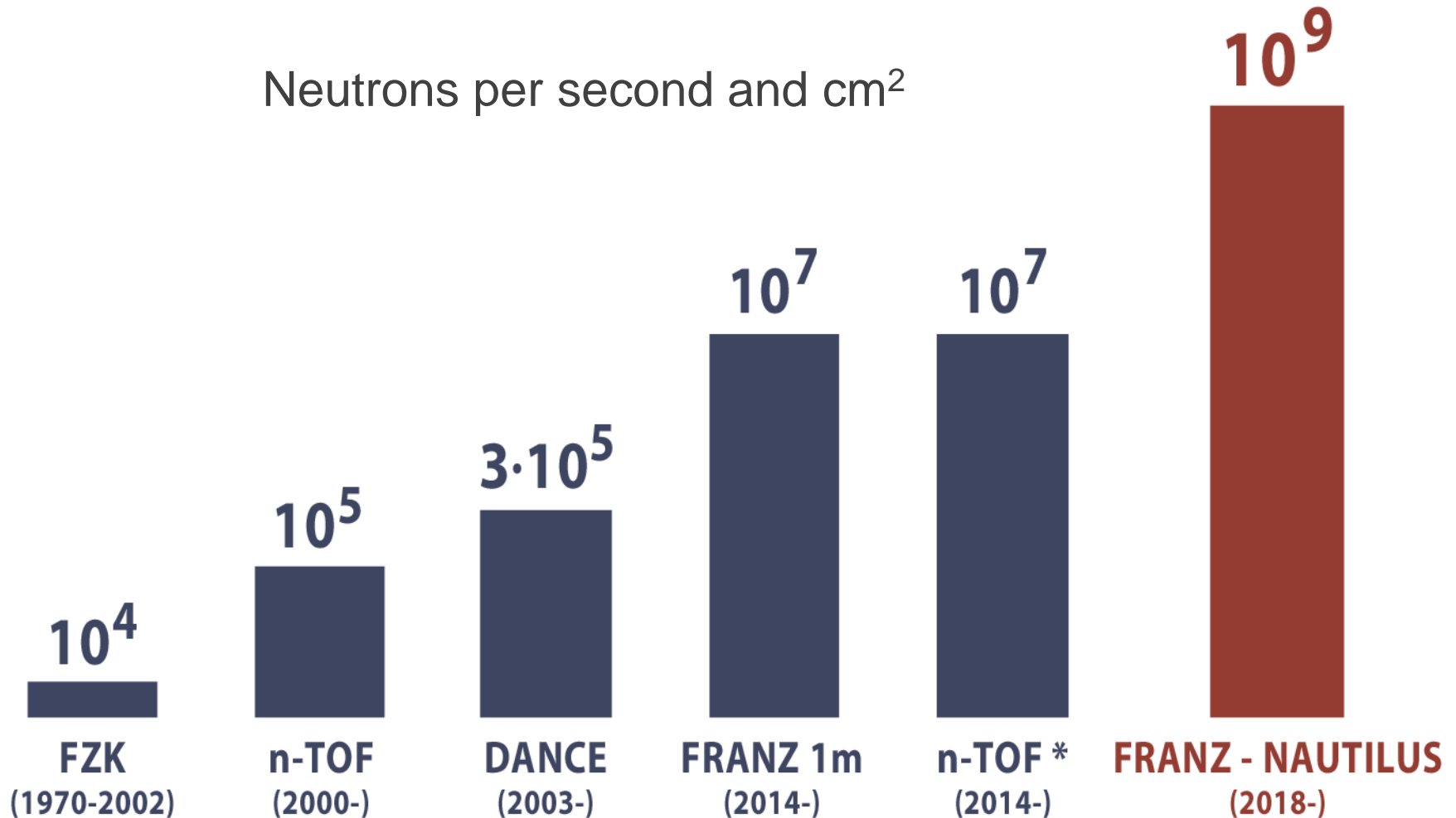
**Increase neutron
flux by factor 100**

NAUTILUS – Expected Time-Of-Flight spectrum



Neutron flux in astrophysical region

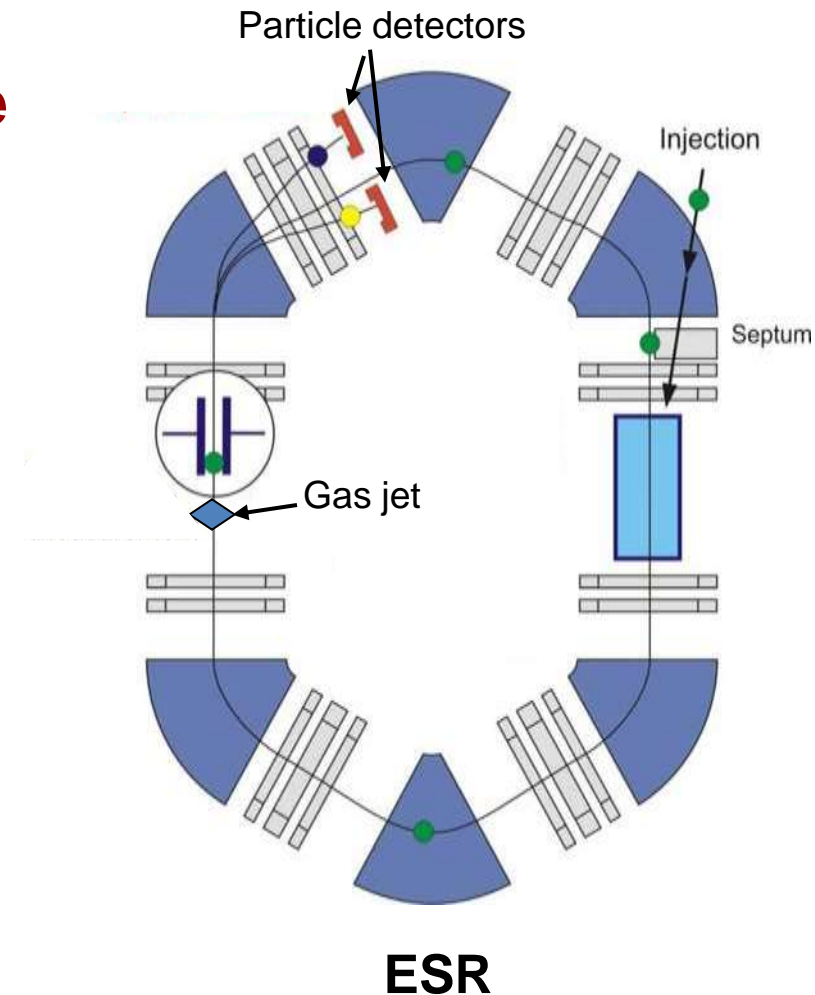
Neutrons per second and cm²



Measurements of (p,γ) or (α,γ) rates in the Gamow window of the p-process in inverse kinematics in the Experimental Storage Ring.

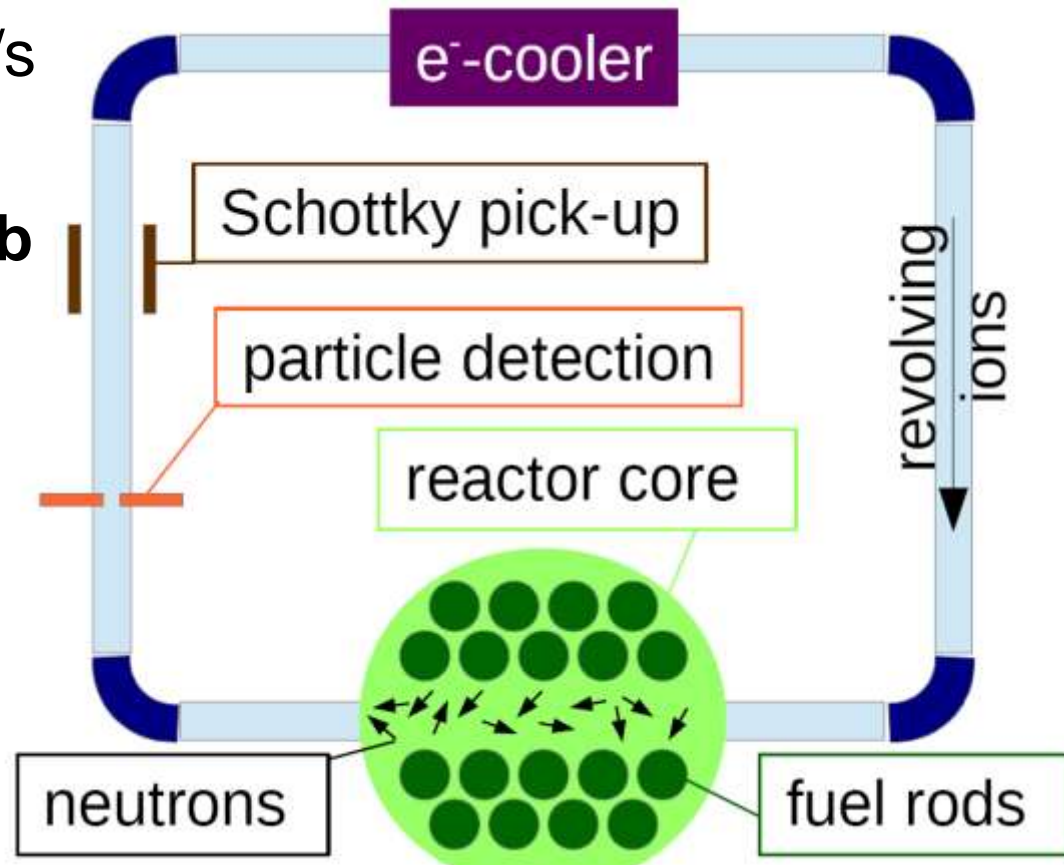
Advantages:

- **Applicable to radioactive nuclei**
- **Detection of ions via in-ring particle detectors (low background, high efficiency)**
- **Knowledge of line intensities of product nucleus not necessary**
- **Applicable to gases**



Neutron captures in inverse kinematics

- Neutron flux: 10^{14} n/cm²/s ->
- Neutron target: $2 \cdot 10^{10}$ n/cm²
- 10^7 ions, 1 MHz: 10^{13} ions/s
- **Counts per day: $20 \sigma / \text{mb}$**



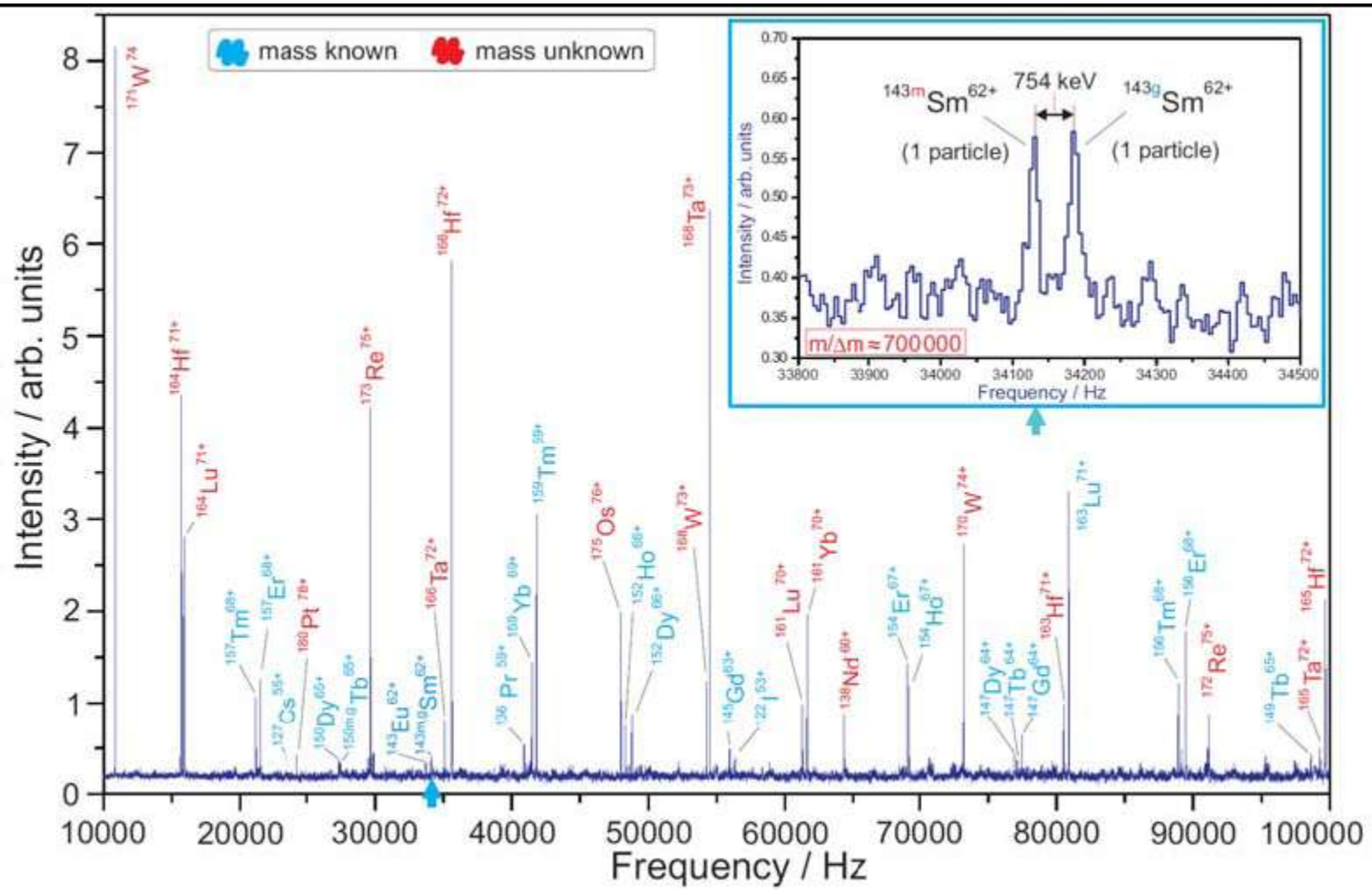
Reifarth & Litvinov, Phys. Rev ST Accelerator and Beams, 17 (2014) 014701

$$B \rho = \frac{mv}{q} = \frac{p}{q} = \text{const}$$

$$\frac{r_{(n,\gamma)}}{r_{\text{primary}}} = 1$$

- Same track as primary beam
- Reacceleration necessary – electron cooler
- Schottky analysis – determine revolution frequency

Schottky Analysis of revolving ions

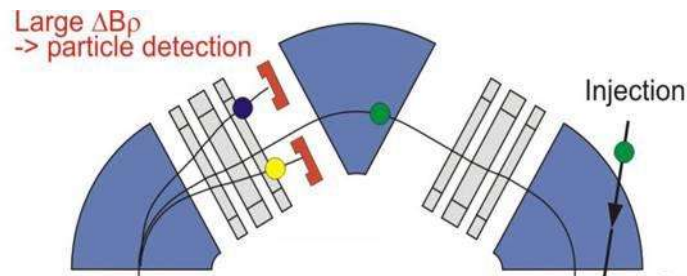


e.g. Y. A. Litvinov and F. Bosch, Rep. Prog. Phys. 74, 016301 (2011)

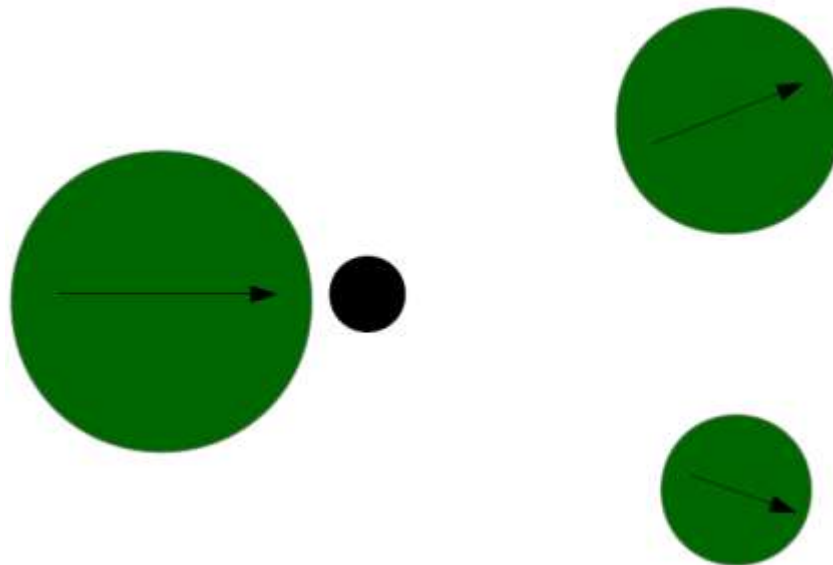
Charged-particle production, (n,2n)

- (n,α) : particle detectors
- (n,p) : particle detectors
- (n,2n) : particle detectors or Schottky

$$\frac{r_{secondary}}{r_{primary}} < 1$$



(n,f) : only at higher energies $E_{\text{CM}} > 10 \text{ MeV}$



- Energy regime: $E_n > 100$ keV
- Half live limit: $t_{1/2} > 0.5$ h, if reactor is not pulsed
- Pulsed reactors might allow even smaller half-lives

- **Radioactive isotopes become more and more in reach of current experimental research**
- **Neutron induced reaction studies are difficult on stable, very difficult on unstable nuclei**
- **FRANZ & NAUTILUS will push the limit further**
- **A combination of a reactor and a ion storage ring might open a new era**